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The energy penalty of post-combustion CO₂ capture & storage and its implications for retrofitting the U.S. installed base

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A review of the literature has found a factor of 4 spread in the estimated values of the energy penalty for post-combustion capture and storage of CO₂ from pulverized-coal (PC) fired power plants. We elucidate the cause of that spread by deriving an analytic relationship for the energy penalty from thermodynamic principles and by identifying which variables are most difficult to constrain. We define the energy penalty for CCS to be the fraction of fuel that must be dedicated to CCS for a fixed quantity of work output. That penalty can manifest itself as either the additional fuel required to maintain a power plant’s output or the loss of output for a constant fuel input. Of the 17 parameters that constitute the energy penalty, only the fraction of available waste heat that is recovered for use and the 2nd-law separation efficiency are poorly constrained. We provide an absolute lower bound for the energy penalty of ~11%, and we demonstrate to what degree increasing the fraction of available-waste-heat recovery can reduce the energy penalty from the higher values reported. It is further argued that an energy penalty of ~40% will be easily achieved while one of ~29% represents a decent target value. Furthermore, we analyze the distribution of PC plants in the U.S. and calculate a distribution for the additional fuel required to operate all these plants with CO₂ capture and storage (CCS).

Introduction

Global carbon dioxide (CO₂) emissions have accelerated from 1.1%/yr in the 1990’s to over 3%/yr since 2000.¹ Those continued growth rates would result in global CO₂ emissions of ~40 GtCO₂/yr and ~100 GtCO₂/yr by 2050, respectively. Stabilizing atmospheric CO₂ concentration below 550 ppm, however, requires emissions to essentially stay flat for the next 42 years.² CO₂ capture and storage (CCS) is a promising technology that has the potential to address the ~40% of emission emanating from large-point sources such as power plants.³ CCS for existing plants involves separating the CO₂ from the plant’s flue gas, compressing the CO₂ for pipeline transport, and injecting the CO₂ into a geologic formation where it is intended to remain for millennia.

The U.S. has 1493 coal-fired power plants that constitute 336 gigawatts (GW) of rated power generation capacity. Nearly all of these plants involve pulverized-coal (PC) combustion, where the coal is pulverized such that over 98% of it is less than 300 μm,⁴ and then it is combusted in air at atmospheric pressure. In 2006, these plants composed about 70% of U.S. fossil-fuel derived electricity and about 50% of total electricity production.⁵ To produce that electricity, the plants burned ~930 million tonnes of coal and produced ~1.9 gigatonnes (Gt) of CO₂, about 1/3 of total U.S. emissions.⁶ In addition, PC power constitutes well over 90% of coal-fired power in the world.⁴ The dominance of PC power plants makes significant reduction in national or global

Broader context

We derive an analytic relationship for the energy penalty from thermodynamic principles, and we apply that relationship to the installed base of U.S. coal-fired power plants to determine the energetic requirements of retrofitting that base for CCS. Pulverized-coal (PC) facilities compose over 95% of the CO₂ produced by U.S. coal-fired power plants. It is unlikely that either national or global CO₂ emissions can be substantially reduced without either shutting down or retrofitting these plants for CCS. The economics of CCS from PC power plants depend, to a large degree, on the thermodynamic work required to capture and store the CO₂. The data demonstrate that—under reasonable assumptions—retrofitting the most efficient 10% of existing plants will offset 30% more CO₂ per unit of additional fuel than retrofitting the least efficient 10% of plants. Indeed, CCS on the least-efficient plants may not make economic sense compared with building new capacity and shutting down the least efficient plants. Finally, we show that a reduction in electrical power demand by between 15% and 20%, combined with retrofitting existing plants for CCS, would lower CO₂ emissions from electricity production by ~65% if the newly liberated power were used for CCS.

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CO₂ emissions from power-plants dependent on either shutting down a substantial fraction of the existing PC plants or retro-fitting those plants for post-combustion capture and storage.

Each of the CCS steps—separation, compression, transport, and storage—requires work. To perform that work, a fraction of the fuel input must be dedicated to CCS. That fuel requirement constitutes the CCS energy penalty. The energy penalty can be realized as either additional fuel input to maintain the baseline power output or as reduced power output for a constant fuel input. Several studies have estimated the energy penalty for PC plants, but these estimates differ between studies by nearly a factor of 4 (Fig. 1).²⁻¹⁸

In this paper, we calculate the thermodynamic work required for the various steps of CCS. We elucidate the cause of the spread in previously estimated energy penalties by deriving an analytic relationship for the energy penalty from first principles and by identifying which of the variables are most difficult to constrain. We show that the energy penalty associated with capturing and storing all the CO₂ generated by U.S. PC plants will require either burning an additional ~400–600 million tonnes of coal per year or building an additional ~100 GW of CO₂-free baseload power. In the latter scenario, the additional baseload power would be required to make up for the reduced power output of the retrofitted PC plants. It should be noted that separate end-use efficiency improvements could serve to offset the CCS energy penalty. In the discussion, we apply our analysis to the actual U.S. fleet of PC power plants. Since the energy penalty is a function of the power-plant’s baseline efficiency, then we derive the expected distribution of energy penalties that would result from retrofitting the U.S. installed base of PC plants.

**Minimum work required to sequester CO₂**

We first derive the lower bound for the work required to capture and store CO₂ from a PC plant. Throughout this section, we neglect frictional losses in both pipeline transport and resistance.
State 4: Concentrated CO₂ stream compressed for injection at
the surface

State 5: Concentrated CO₂ stream emplaced beneath the pore
water in the geologic formation

The process of sequestering CO₂ requires the input of work to
transfer the system from state 1 to states 2 and 3, from state 3 to
state 4, and from state 4 to state 5. We have labeled the work
required for these three transitions: \( W_a \), \( W_b \), and \( W_c \).

\( W_a \) (from state 1 to states 2 and 3)

Separate the CO₂ from the flue gas. Separating the CO₂ from
flue gas is justified because separating and venting gases other than
CO₂ back to the atmosphere reduces the net sequestration work
by lowering the compression and injection costs. In the discussion,
we calculate the optimal separation fraction that minimizes the
total primary energy requirement for a range of efficiencies.

Flue gas emitted from typical coal-burning power plants
contain \( \sim 78\% \) N₂ from the atmosphere, \( \sim 15\% \) CO₂ from the
oxidation of the carbon in the hydrocarbon, and \( \sim 7\% \) water from
both the oxidation of hydrogen in the coal and the vaporization of
water that was adsorbed on the coal. Removing H₂O from the flue
gas is—in principle—thermodynamically favorable because H₂O
condenses at surface conditions indicating the enthalpy change of
separation is not zero. In the thermodynamic limit, the minimum
work to separate the flue gas into one concentrated CO₂ stream
and one concentrated N₂ stream is the difference in the thermo-
dynamic availability before and after separation. For an
isothermal and isobaric process, the work equals the change in
free energy before and after separation:

\[
W_{\text{min}} = -dG
\]  

(1)

where \( G \) is the Gibbs free energy.

For the separation, we employ the ideal gas assumption as the
pressure is near atmospheric and N₂ and CO₂ do not chemically
interact. The mole fraction of N₂ and CO₂ in the fully mixed state
is \( X_{\text{N}} \) and \( X_{\text{C}} \). After the separation, the mole fraction in state 2
(low CO₂) are \( X_{\text{N2}} \) and \( X_{\text{C2}} \), and the mole fractions for state 3
(high CO₂) are \( X_{\text{N3}} \) and \( X_{\text{C3}} \).

The partial molar Gibbs energy for each gas in an ideal
mixture is given by:

\[
\frac{\partial G}{\partial n_i} = G_i^0 + RT \ln \left( \frac{P_i}{P} \right)
\]

where \( P_i \) is the partial pressure of the \( i \)th gas and \( P \) is the total
pressure. Thus, the total free energy of an ideal gas mixture is:

\[
G_{\text{tot}} = \sum_i n_i \frac{\partial G}{\partial n_i}
\]

(3)

If we assume that none of the states are completely pure \((i.e.,
X_{ij} > 0 \text{ for all } i \text{ and } j)\), then the previous equation will determine

the water table is independent of the size of the domain and the geometry
of the injected CO₂ plume. If CO₂ were injected beneath twice the land
area, then the change in potential energy would not change because
a greater quantity of water would be lifted a corresponding smaller
distance.
the free energy of each state. The minimum work to change a system’s state is given by the change in free energy between those states:

$$W = \Delta G_{\text{cap}} = (G_2 + G_3) - G_1$$

To calculate the minimum work required to transfer the system from state 1 into the distinct states 2 and 3, we calculate the free energy of each state:

$$G_1 = n_N G_0 + n_{CO_2} G_0^b + RT (n_{CO_2} \ln(X_{CO_2}) + n_{N_2} \ln(X_{N_2}))$$
$$G_2 = n_{CO_2} G_0^f + n_{N_2} G_0^b + RT (n_{CO_2} \ln(X_{CO_2}) + n_{N_2} \ln(X_{N_2}))$$
$$G_3 = n_N G_0 + n_{N_2} G_0^b + RT (n_{N_2} \ln(X_{N_2}) + n_{N_1} \ln(X_{N_1}))$$

So, the minimum required work is:

$$W = (RT(n_{CO_2} \ln(X_{CO_2}) + n_{N_2} \ln(X_{N_2}))) + RT(n_{N_2} \ln(X_{N_2}) + n_{N_1} \ln(X_{N_1})) - RT(n_{CO_2} \ln(X_{CO_2}) + n_{N_2} \ln(X_{N_2}))$$

$$\text{(6)}$$

And the work per mole of CO$_2$ in the flue gas is:

$$W = \frac{1}{n_{CO_2}} (RT(n_{CO_2} \ln(X_{CO_2}) + n_{N_2} \ln(X_{N_2}))) + RT(n_{N_2} \ln(X_{N_2}) + n_{N_1} \ln(X_{N_1})) - RT(n_{CO_2} \ln(X_{CO_2}) + n_{N_2} \ln(X_{N_2}))$$

$$\text{(7)}$$

The number of parameters can be reduced by substituting the definition of the mole fraction into eqn (6), and the minimum work required to isothermally separate an ideal gas mixture into two ideal gas mixtures of different concentrations ($W_a$) is given by:

$$W_a = RT \left[ \begin{array}{c} n_{CO_2} \ln \left( \frac{n_{CO_2}}{n_{CO_2} + n_{N_2}} \right) + n_{N_2} \ln \left( \frac{n_{N_2}}{n_{CO_2} + n_{N_2}} \right) \\
+ (n_{N_1} - n_{CO_2}) \ln \left( \frac{n_{CO_2}}{n_{CO_2} + n_{N_2}} \right) \\
+ (n_{N_1} - n_{CO_2}) \ln \left( \frac{n_{N_1} - n_{CO_2}}{n_{CO_2} + n_{N_2} - n_{CO_2}} \right) \\
- \left( n_{CO_2} \ln \left( \frac{n_{CO_2}}{n_{CO_2} + n_{N_2}} \right) + n_{N_2} \ln \left( \frac{n_{N_2}}{n_{CO_2} + n_{N_2}} \right) \right) \end{array} \right]$$

$$\text{(8)}$$

where $n$ is the number of moles of either N$_2$ or CO$_2$, subscripted N or C, in either the original mixture, the concentrated CO$_2$ stream, or the concentrated N$_2$ stream, subscripted 1, 2, and 3 respectively. For typical values, $W_a$ equals ~9 kJ/(mol CO$_2$).

$W_b$ (from state 3 to state 4)

Compress the concentrated stream at the surface. To inject the concentrated CO$_2$ stream into a geologic formation, it must be compressed such that at the bottom of the well, its pressure equals the reservoir pore-pressure:

$$P_b = P_s - g \int_0^{L_4} \rho_s(z) dz = \rho_w g L_d - g \int_0^{L_4} \rho_s(z) dz$$

$$\text{(9)}$$

$P_b$ is the pressure at the top of the well, where subscript 4 indicates the state of concentrated and compressed CO$_2$ at the surface. $P_s$ is the pore-pressure at the bottom of the well, which is assumed to initially equal the product of the density of water ($\rho_w$), the gravitational constant ($g$), and the depth of injection ($L_d$). The integral in the second term accounts for gravitational compression and thermal expansion of the concentrated CO$_2$ stream within the borehole.

The minimum work required to compress the concentrated CO$_2$ stream from state 3 to state 4 ($W_b$) is given by the reversible isothermal compression:

$$W_b = - \int_{\gamma = \gamma_1(T_3, T_4)}^{\gamma = \gamma_2(T_3, T_4)} (P(v, T) - P_a) dv$$

$$\text{(10)}$$

where $v_1$ is the molar volume of the concentrated CO$_2$ stream, $P_3$ and $T_3$ are the conditions at which the plant supplies the highly concentrated stream, $P_4$ and $T_4$ are the post compression conditions, and $P_a$ is the atmospheric pressure, which assists in the compression. For typical values, $W_b$ equals ~13 kJ/(mol CO$_2$).

$W_c$ (from state 4 to state 5)

Push the compressed CO$_2$ into the formation. Once the CO$_2$ at the surface is compressed to the pressure at the bottom of the borehole ($P_b$), it must be pushed out of the compressor and into the well, which causes CO$_2$ at the well screen to flow into the formation and vertically displace the ground-water. In the limit of zero friction, pushing the CO$_2$ into the formation requires work to vertically displace the ground-water ($W_c$) and work to create the interfacial surface between the CO$_2$ and the pore water ($W_{c2}$).

The work required to vertically displace the ground water is equal to the volume of ground water displaced times the well-head pressure:

$$W_{c1} = (P_4 - P_a) v_4$$

$$\text{(11)}$$

In the limiting case of zero viscous drag, the minimum value of $W_{c1}$ occurs when $P_4$ equals the hydrostatic pressure of the ground water at the well-head. For typical values, $W_{c1}$ equals ~1-2 kJ/(mol CO$_2$).

In addition, work is required to create the interfacial surface between the CO$_2$ and the pore-water because, on the injection time-scale, the CO$_2$ acts primarily as an immiscible phase.$^{21}$ The interfacial surface tension between supercritical CO$_2$ and water at the relevant conditions is ~0.02 J/m$^2$,$^{22}$ and the work required to increase the surface area goes as the interfacial surface tension ($\gamma$) and the change in surface area of the interface:

$$W_{c2} = \gamma dA = \Delta V \Delta P_{\text{cap}}$$

$$\text{(12)}$$

$W_{c2}$ is the work required to overcome the capillary force; where $\Delta V$ is the total swept out pore volume, and $\Delta P_{\text{cap}}$ is the capillary pressure:$^{23}$

$$\Delta P_{\text{cap}} = \gamma \left( \frac{1}{R_1} + \frac{1}{R_2} \right)$$

$$\text{(13)}$$

$R_1$ and $R_2$ are the radii of curvature of the surface, and for a typical sand-reservoir, the pore sizes are on the order of
10^{-4} \text{--} 10^{-5} \text{m}^3. The molar volume of CO$_2$ at reservoir conditions is \( \sim 10^{-4} \text{m}^3 \). Therefore, the capillary pressure is on the order of 2000 Pa, and \( W_{c2} \) is on the order of 0.2 J/(mol CO$_2$). Since \( W_{c1} \), on the other hand, is on the order of \( \sim 10^4 \text{J/(mol CO}_2) \), then \( W_{c2} \) can be safely ignored.

The total compression work is the sum of \( W_b \) and \( W_c \), where \( W_c \) is the sum of \( W_{c1} \) and \( W_{c2} \). By the reverse integration by parts:

\[
W_{bc} = W_b + W_c = - \int_{P_4}^{P_3} (P-v) \, dp + (P_4-P_a) v_4
\]

W$_{bc}$ is minimized under isothermal conditions, but regardless of the final temperature, the minimum work input required to achieve this lower bound injection pressure equals the change in thermodynamic availability of the concentrated CO$_2$ stream during compression:

\[
W_{bc} = n_C((h_C4 - h_C3) - T_0(s_C4 - s_C3)) + n_N((h_N4 - h_N3) - T_0(s_N4 - s_N3))
\]

where \( h \) and \( s \) are the molar enthalpies and entropies, respectively. The subscripts, 3 and 4 indicate the surface conditions before and after compression, and \( T_0 \) is the ground state temperature. Integration of eqn (14) from \( P_3 \) to \( P_4 \) will always equal the state change of the thermodynamic availability (eqn (15)) as shown in Fig. 3.

A different approach to ascertain the work required to emplace the CO$_2$ below the pore water is to calculate the change in gravitational potential energy between state 4 and state 5. The total turbine shaft work is then the sum of \( W_b \) and the change in gravitational potential energy:

\[
W_{bc} = W_b + \Delta U = W_b + g L \frac{\rho_b}{\rho_{C5}} \left( m_C/n_C + m_N/n_N \right)
\]

where \( m_C \) and \( m_N \) are the molar masses of CO$_2$ and N$_2$ respectively, \( \rho_i \) are densities of H$_2$O, CO$_2$, and N$_2$ in state 5, and \( \Delta U \) is the change gravitational potential energy of the system from state 4 to state 5.

Eqn (14), (15), and (16) all yield the same value for \( W_{bc} \), the minimum possible work. It should be noted that this lower limit depends on the depth of injection, but it is independent of the reservoir geometry and area [see Fig. 2(c)]. That is to say that creating sufficient pore-space for a given volume of CO$_2$ in a reservoir with smaller area requires greater vertical displacement of less pore-water while emplacing the CO$_2$ into a reservoir with larger area requires less vertical displacement of more pore-water. The displaced volume and the required work are the same in both cases.

Friction. The energetic calculations are now extended to include friction associated with transport and storage. As the CO$_2$ is transported from the injection zone to the storage site, viscous drag within the pipeline decreases the fluid pressure. The pressure drop in a pipeline is given by:

\[
\Delta P = \rho \beta \left( \frac{LV^2}{2D} \right)
\]

where \( \rho \) is the fluid density, \( L \) is pipeline length, \( V \) is the fluid velocity, \( D \) is the pipeline diameter, and \( \beta \) is the friction factor, which depend on the Reynolds number of the flow and inner roughness of the pipeline.\( ^{26} \) For typical values \( (i.e., \rho = 900 \text{ kg/m}^3, D = 0.5 \text{ m}, V = 2 \text{ m/s} \beta \sim 10^{-2} \text{ and } \Delta P \sim 50 \text{ L (kg/m}^2\text{s}^2) \). Therefore, a typical pressure drop is \( \sim 0.5 \text{ bars per km or 25 bars per 50 km} \). Eqn (14) indicates that—in the thermodynamic limit—the work required to re-compress the CO$_2$ from 90 bars to 115 is \( \sim 0.4 \text{ kJ/(mol) per 50 km} \).

Friction is encountered again when CO$_2$ is injected into the formation and work is required to compress the pore-fluid and expand the pore-space. That work dissipates into friction, which is manifest in elevated injection pressure requiring additional compression work \( (W_b \text{ and } W_c) \). The frictional loss is larger for high rates of flow—and so the compression work is greater for faster injection. But, storage within the reservoir reduces work by slowing down flow.

Induced fracturing, however, limits reservoir pressures to the formation’s least compressive stress.\( ^{27} \) The fracture threshold is predominately determined by the least compressive stress. If the vertical stress, which is often near lithostatic pressure, is the weakest stress, then horizontal fractures will form and accommodate the injected CO$_2$. It is more likely, however, that one of the horizontal stresses is the weakest indicating that vertical fractures are most likely.\( ^{28} \)

In any case, the pressure and the required work are bounded by pressure-induced fracturing. Prior to having reached fracture

![Fig. 3](https://example.com/image3.png)

The left-hand vertical axis is for the volume–pressure curves of an ideal gas and pure CO$_2$ at 340 K. The non-ideal behavior of CO$_2$ at above \( \sim 70 \text{ bars indicates that less work is required to fully compress CO}_2 \text{ than would be required to compress an ideal gas. The right-hand vertical axis depicts the thermodynamic availability of a pure CO}_2 \text{ phase as a function of pressure at 340 K. The definite integral of the volume–pressure curve between two pressures equals the state change in thermodynamic availability between those same two pressures.}
pressure, however, additional compression work is required to maintain the constant flow of CO₂ into the formation. Over the lifetime of the injection, the reservoir pressures can increase from the hydrostatic pressure of the formation to the fracture pressure, which is ~1.4–2.0 times the hydrostatic pressure.²⁷

The additional compression work required (i.e., the increase in \( W_b \)) to accommodate the increased reservoir pressure is small compared with the initial compression as the CO₂ pressure-volume curve (Fig. 3) flattens out significantly above 100 bars. On the other hand, the additional work required to emplace the CO₂ beneath the pore-water by pushing the fluid into the reservoir (i.e., the increase in \( W_c \)) is on the same order as the initial work \( W_c \). That is because eqn (11) is linear in pressure, and the fracture pressure is about twice the initial hydrostatic pressure. Specifically, \( W_b \) changes from ~12 kJ/(mol CO₂) initially to 13 kJ/(mol CO₂) once fracture pressure is reached, while \( W_c \) changes from 1 kJ/mol to ~2 kJ/mol. Thus, the friction related work increases \( W_{bc} \) by ~15%.

**Primary energy required and waste heat recovery**

The minimum total work required to separate, compress, transport, and store a unit of CO₂ produced from a power plant is:

\[
W_{\text{tot}} = W_a + W_b + W_c
\]  

(18)

Performing work \( W_a \), \( W_b \), and \( W_c \), requires primary energy \( E_a \), \( E_b \), and \( E_c \). Calculating \( E_a \)—the sum of \( E_a \), \( E_b \), and \( E_c \)—requires determining the theoretical limit for each conversion of primary energy to work. Subsequently, we extend our calculations to consider reasonable upper bounds on efficiencies, and hence reasonable lower bounds for the energy penalty.

\( E_a \)

A well developed technology for post-combustion separation of CO₂ from flue gas is the temperature-swing system with the solvent monoethanolamine (MEA).²⁹ \( E_a \) is not set by the enthalpy of desorption because nearly the same quantity of heat is generated during the exothermic absorption as the endothermic desorption. Rather, \( E_a \) is set by the 2nd-law of thermodynamics, which limits the efficiency of any heat-to-work conversion.²⁹ Therefore, \( E_a \) is given by \( W_a \) over the product of the ideal heat-to-work conversion efficiency for separation (\( \eta_{\text{ideal}} \)) and the 2nd-law efficiency of the actual separation process (\( \eta_{\text{act}} \)):

\[
E_a = \frac{W_a}{\eta_{\text{ideal}} \eta_{\text{act}}} \quad \text{n} \quad (19)
\]

\( \eta_{\text{act}} \) for a temperature-swing separation—such as monoethanolamine (MEA)—is given by the efficiency of a Carnot engine running between its highest and lowest temperatures. The lowest temperature (\( T_a \sim 310 \text{ K} \)) of an MEA system is the temperature of the absorber unit that binds lean MEA to CO₂, while the highest temperature (\( T_S \sim 390 \text{ K} \)) is the temperature of the of the stripper unit that desorbs the rich MEA (Fig. 2);²⁹

\[
\eta_{\text{act}} = 1 - \frac{T_S}{T_a} \quad \text{n} \quad (20)
\]

Temperature-swing separation systems—like MEA—often work at low temperatures.²⁹ For that reason, harnessing low-temperature waste heat for this type of separation is practical, and as such, the additional primary energy required for separation, \( E_{\text{sep}} \), is lower than \( E_a \). If the quantity of waste heat that can—in principle—be used for separation is \( E_w \), then and the primary energy required from additional combustion of fuel is:

\[
E_{\text{sep}} = \frac{W_a}{\eta_{\text{ideal}} \eta_{\text{act}}} - E_w \quad \text{n} \quad (21)
\]

Or,

\[
E_{\text{sep}} = \frac{W_a}{(1 - \frac{T_a}{T_s}) \eta_{\text{act}}} - E_w \quad \text{n} \quad (22)
\]

\( E_w \)—the available waste heat—is calculated by considering the Carnot efficiency of the power plant:

\[
\eta_c = 1 - \frac{T_L}{T_H} = 1 - \frac{E_L}{E_H} \quad \text{n} \quad (23)
\]

where \( T_L \) is the minimum temperature of the working fluid (i.e., the condenser), \( T_H \) is maximum temperature of the working fluid (i.e., modern-ultra-critical steam cycles run at ~1000 K⁴), \( E_L \) is the minimum quantity of heat transferred to the environment, and \( E_H \) is the primary energy content of the fuel, which for a high rank coal is ~400 kJ/(mol CO₂).²⁰

The actual power plant efficiency (\( \eta_{\text{pp}} \)) is given by:

\[
\eta_{\text{pp}} = 1 - \frac{E_L}{E_H} \quad \text{n} \quad (24)
\]

where \( E_L \) is the actual quantity of heat dumped to the environment, which is equal to \( E_L \) plus the total waste heat produced by irreversible processes (or inefficiencies) within the plant (\( E'_w \)). Thus,

\[
\eta_{\text{pp}} = 1 - \frac{E_L + E'_w}{E_H} \quad \text{n} \quad (25)
\]

Solving for \( E'_w \) with the above 3 equations yields:

\[
E'_w = E_H(\eta_c - \eta_{\text{pp}}) \quad \text{n} \quad (26)
\]

Eqn (26) gives the total waste heat produced by irreversible processes (or inefficiencies) within the plant by a power plant. \( E'_w \), however, is greater than \( E_w \)—the available waste—because temperature-swing separation systems have a minimum temperature at which heat can be used. The stripper accepts heat at or above \( T_S \) while the absorber dumps heat to the environment (or the condenser) at or above \( T_L \) (Fig. 2). As such, waste heat can only be utilized if it is greater than \( T_S \). The fraction of available waste heat (\( \eta_w \)) approaches zero as \( T_S \) approaches \( T_H \). Likewise, \( \eta_w \) approaches 1 as \( T_S \) approaches \( T_L \). There is limited published data on the temperature distribution of power-plant waste heat. For that reason, a linear temperature distribution between \( T_L \) and \( T_H \) was assumed. Such a waste–heat distribution yields:

\[
\eta_w = \frac{T_H - T_S}{T_H - T_L} \quad \text{n} \quad (27)
\]
Careful research is needed to more accurately access the waste–heat distributions. Such work would be very valuable, but also very difficult as it would require detailed engineering studies of existing power-plants. The goal of this study is to constrain what is physically possible given actual and theoretical power plant efficiencies.

Given our assumed waste–heat distribution, the quantity of available waste heat is:

\[ E_w = E_{\text{th}} \left( \eta_c - \eta_{pp} \right) \frac{T_H - T_S}{T_H - T_L} \]  

(28)

Or,

\[ E_w = E_{\text{th}} \left( \left( 1 - \frac{T_L}{T_H} \right) - \eta_{pp} \right) \frac{T_H - T_S}{T_H - T_L} \]  

(29)

Only a fraction of the available waste heat, however, will actually be harnessed. Therefore, we introduce a new variable \( \eta_{w_{2nd}} \), which is the 2nd-law efficiency (i.e., the fraction of maximum possible) of the waste heat recovery, and therefore, the total quantity of waste heat that is actually recovered for productive use is:

\[ E_{w_{2nd}} = E_{\text{th}} \left( \left( 1 - \frac{T_L}{T_H} \right) - \eta_{pp} \right) \frac{T_H - T_S}{T_H - T_L} \eta_{w_{2nd}} \]  

(30)

Combining eqn (21) and (30) yields the additional primary energy required to perform \( W_a \) for a temperature-swing separation process:

\[ E_{\text{sep}} = E_1 - E_{w_{2nd}} \]

\[ = \frac{W_a}{1 - \frac{T_S}{T_H}} \eta_{2nd} - E_{\text{th}} \left( \left( 1 - \frac{T_L}{T_H} \right) - \eta_{pp} \right) \frac{T_H - T_S}{T_H - T_L} \eta_{w_{2nd}} \]  

(31)

**Energy and Efficiency**

The primary energy required for compression depends on the efficiency of the turbine to produce shaft work and the 2nd-law efficiency of the compressors themselves. No waste heat is allocated to compression as that would require building a more efficient power plant. Therefore:

\[ E_b + E_c = \frac{W_b + W_c}{\eta_{pp} \eta_{com}} \]  

(32)

where \( \eta_{pp} \) is the power plant thermal efficiency and \( \eta_{com} \) is the isothermal compressor efficiency. Available compressors report isothermal efficiencies of between \( \sim 57\% \) and \( 66\% \).

The re-compression necessary to overcome pipeline and reservoir friction will be powered by either natural-gas from a parallel pipeline or electricity from the grid. Calculating the associated CO2 emissions requires either the efficiency of natural gas compressors or the collective efficiency of grid power with respect to CO2 emissions (i.e., all the electric power produced over all the thermal power from CO2 production). Assuming the CO2 is not transported more than a few hundred kilometers, the primary energy calculated from such values increases \( E_{bc} \) by less than 10%.

Adding eqn (31) and (32), we arrive at the lower limit of the total primary energy used by a temperature-swing separation system for CO2 sequestration:

\[ E_s = \frac{W_a}{1 - \frac{T_S}{T_H}} \eta_{2nd} - E_{\text{th}} \left( \left( 1 - \frac{T_L}{T_H} \right) - \eta_{pp} \right) \frac{T_H - T_S}{T_H - T_L} \eta_{w_{2nd}} \]  

\[ + \frac{W_b + W_c}{\eta_{pp} \eta_{com}} \]  

(33)

For non-temperature-swing separation systems—such as membranes or pressure-swing absorption systems—waste heat is not useful because the separation work is a parasitic load on the power-plant turbine. As such, \( E_s \) is:

\[ E_s = \frac{W_a}{\eta_{2nd} \eta_{pp}} + \frac{W_b + W_c}{\eta_{com} \eta_{pp}} \]  

(34)

**The energy penalty**

The energy penalty, \( f_1 \), is the fraction of the fuel that must be dedicated to CCS activities for a given quantity of fuel input (\( E_{\text{th}} \)):

\[ f_1 = \frac{E_s}{E_{\text{th}}} \]  

(35)

If the quantity of fuel is fixed, then the energy penalty is manifest in the reduction of the plant’s power output. The new power output per unit fuel is:

\[ W' = \eta_{pp} E_{\text{th}} \left( 1 - \frac{E_s}{E_{\text{th}}} \right) = \eta_{pp} E_{\text{th}} (1 - f_1) \]  

(36)

If, on the other hand, the power output is fixed, then the energy penalty is manifest as the increase in fuel necessary to maintain that constant power output. This additional fuel requirement is expressed as the ratio of the fuel for CCS (\( E_s \)) to the fuel that produces power output (\( E_{\text{th}} - E_s \)). Thus, the fraction of additional fuel required to maintain the constant power output associated with \( E_{\text{th}} \) is:

\[ f_2 = \frac{E_s}{E_{\text{th}} - E_s} = \frac{E_s/E_{\text{th}}}{E_{\text{th}}/E_{\text{th}} - E_s/E_{\text{th}}} = \frac{f_1}{1 - f_1} \]  

(37)

This additional fuel requirement can also derived as the geometric sum of the energy penalty. The additional fuel required to sequester the CO2 produced for a unit of power generation is \( f_1 \), but in burning this fuel, more CO2 is produced that requires \( (f_1)^2 \) additional fuel to sequester, which in turn requires more fuel, ad infinitum, to give the sum of the infinite geometric series:

\[ f_2 = f_1 + (f_1)^2 + (f_1)^3 + ... = \frac{f_1}{1 - f_1} \]  

(38)

By combining eqn (33) with eqn (35), the energy penalty (\( f_1 \)) and the additional fuel requirement (\( f_2 \)) for a temperature-swing separation system can be written in terms of basic system parameters as well as the minimum work required for separation and compression:
\begin{align*}
f_1 &= \frac{E_S}{E_H} = \frac{1}{E_H} \left[ \frac{W_s}{1 - \frac{T_s}{T_H}} - E_H \left( 1 - \frac{T_L}{T_H} \right) - \eta_{pp} \right] \\
&= \frac{T_H - T_S}{T_H - T_L} \frac{T_H - T_S}{T_H - T_L} \eta_{w2nd} + W_b + W_c \eta_{pp} \eta_{com} \\
&= \left( 1 - \frac{T_S}{T_H} \right) \eta_{w2nd} \end{align*}

(39)

and

\begin{align*}
f_2 &= \frac{E_S}{E_H} - \frac{E_S}{E_H} \\
&= \frac{W_s}{1 - \frac{T_S}{T_S}} - E_H \left( 1 - \frac{T_L}{T_H} \right) \eta_{pp} + \frac{W_b + W_c}{\eta_{pp} \eta_{com}} \\
&= \left( 1 - \frac{T_S}{T_H} \right) \eta_{w2nd} \\
&= \left( 1 - \frac{T_S}{T_H} \right) \eta_{w2nd} + W_b + W_c \eta_{pp} \eta_{com} \\
&= \left( 1 - \frac{T_S}{T_H} \right) \eta_{w2nd} \\
&= \left( 1 - \frac{T_S}{T_H} \right) \eta_{w2nd} \end{align*}

(40)

Discussion

The lower bound of the total CCS work—for a 2 km injection—with perfect 2nd-law efficiencies for all three steps is $W_{abc} \approx 24$ kJ/(mol of CO$_2$), where $W_a$, $W_b$, and $W_c$ equal 0.9, 13, and 2 kJ/(mol of CO$_2$). Eqn (33) and (34) convert the work lower bound to the primary energy lower bound for temperature-swing separation systems and pressure-swing separation systems, respectively. For a pressure-swing separation process with perfect compression and perfect separation, the ideal primary energy requirement for the total installed base would be $\sim 75$ kJ/(mol of CO$_2$), which implies $f_1$ of $\sim 19\%$ and $f_2$ of $23\%$ (assuming $E_H \sim 400$ kJ/(mol of CO$_2$)). For a temperature-swing separation process, the minimum primary energy requirement corresponds to the case when sufficient waste heat is harnessed for complete separation. In such an ideal case, the minimum energy requirement would be $\sim 45$ kJ/(mol CO$_2$), implying an $f_1$ of $\sim 11\%$ and an $f_2$ of $\sim 13\%$.

Table 1 reveals the end-member cases for post-combustion capture and storage between the thermodynamic lower-bound and values being reported for current technology. From these values, it is clear that the energy penalties achieved from temperature-swing separation systems are more uncertain, but also that waste–heat recovery offers a significant opportunity to decrease the energy penalty.

The thermodynamic limit for sequestration with a temperature-swing separation system indicates that capturing and storing all the CO$_2$ generated from current U.S. PC plants while delivering the same power output would require—at the very least—considering an additional $\sim 120$ million tonnes of coal annually. If, on the other hand, the energy penalty were incurred by decreasing the electrical work output—rather than increasing the fuel consumption—then the electrical work output of the U.S. coal fleet will drop by—at the very least—$\sim 37$ GW. That means that either an additional 37 GW of base-load CO$_2$-free power have to be built, or national electricity use would have to be reduced by 37 GW.

The energy penalty for post-combustion capture and storage of PC power-plant CO$_2$ has been estimated by several different studies. A review of that literature demonstrates a relationship between the energy penalty and the economics of CCS (Fig. 1); it also reveals a significant spread between the various published estimates of the energy penalty. The reviewed studies of PC retrofits include only small amounts of waste–heat recovery, and the associated $f_2$ values are between 43% and 77%, which indicate $\eta_{w2nd}$ values of $\sim 40\%$–60% (Fig. 1). Analysis of eqn (39) and (40) provides insight into that spread as well as into future CCS development.

Eqn (40) depends on 17 parameters, yet in practice only two of them—the fraction of available waste heat that is actually recovered ($\eta_{w2nd}$) and the 2nd-law separation efficiency ($\eta_{com}$)—are poorly constrained ($\eta_{pp}$ varies significantly, but it is well constrained). Fig. 4 reveals how $f_2$ depends on these poorly constrained parameters.

New construction projects have two distinct advantages; first, the power plants themselves—having supercritical or ultracritical steam cycles—are more efficient which results in a lower energy penalty since the primary energy required for compression is function of $\eta_{pp}$, and second, new construction projects can more easily be designed to utilize low-grade waste heat for CO$_2$ separation. The absolute value of the contour slopes in Fig. 4 are mostly less than 1, which indicates that $f_2$ is generally more sensitive to changes of $\eta_{w2nd}$ than to changes of $\eta_{pp}$. Since current systems already achieve $\eta_{w2nd}$ values in the range of $\sim 40\%$–60%($\eta_{w2nd}$ represents the most potential for
improvement as the contour slopes are significantly steeper for $h_{s2nd}$ values above 50%.

The energy penalty—as derived in this paper—can be used to calculate the optimal values for various independent variables. For instance, $W_a$ and $W_{bc}$ depend on the degree of separation in opposite directions. Zero separation minimizes $W_a$, but it maximizes $W_{bc}$. The same is true—naturally—for $E_a$ and $E_{bc}$.

Fig. 5 shows the optimal separation as a function of the fraction of flue-gas CO$_2$ that is emitted to the atmosphere.

The stripper temperature is another parameter for optimization as $T_S$ affects $f_1$ in two different directions. The ideal separation efficiency increases with $T_S$ (eqn (20)), but the quantity of available waste heat decreases with $T_S$ (eqn (27)). MEA strippers operate at $\sim$390 K, and research is ongoing to identify new absorption materials—such as ionic liquids$^{32}$—that can operate at higher temperatures with the goal of increasing the heat to work conversion efficiency. Due to the decrease in available waste heat, however, these efforts might be limited in their potential. For typical efficiency and waste–heat recovery values,
increasing the $T_s$ beyond $\sim 500$ K might not be helpful because the loss of available waste heat compensates for the increase in separation efficiency [Fig. 6(a)]. Indeed, given the likelihood that the waste–heat temperature distribution is more skewed toward lower temperatures than the linear distribution assumed here, it is probable that the optimal $T_s$ is below 500 K, suggesting that current systems are operating near the optimal $T_s$. That assumes, however, that effective engineering can harness the available waste heat.

Sensitivity analysis on the total energy penalty is performed by varying the power-plant efficiency. The additional fuel requirement ($f_2$) depends on the power-plant efficiency in two ways: $E_{\text{bc}}$ decreases as $\eta_{\text{pp}}$ increases, but $E_w$ can actually decrease as $\eta_{\text{pp}}$ increases because the available waste heat decreases as $\eta_{\text{pp}}$ increases. Fig. 6(b) shows $f_2$ as a function of $\eta_{\text{pp}}$, and that figure demonstrates that $f_2$ is minimized for particular power-plant efficiencies. If available waste heat recovery rates can exceed 30%, then it may not be beneficial to target inefficient plants for CCS retrofits.

The reviewed studies also indicate significant differences in the energy penalty between new construction projects and retrofits. Those differences are primarily driven by 3 factors that are made clear from our analysis of the energy penalty: (1) the degree of available-waste-heat recovery ($\eta_w$), (2) the baseline power plant efficiency ($\eta_{\text{pp}}$), and (3) the 2nd-law separation efficiency ($\eta_{\text{sep}}$). All the studies of new construction projects involve either supercritical or ultra-supercritical cycles whose superior plant efficiencies result in lower energy penalties than subcritical cycles. In addition, waste–heat recovery for separation is easier to implement in new construction projects than in retrofits.

The U.S. installed base of PC plants has a total thermal efficiency ($\eta_{\text{pp}}$) of 33%. Fig. 7(a) shows the distribution of thermal efficiencies for the installed base of PC plants.13

In 2007, the most efficient plant recorded a thermal efficiency of 46.4% while the least efficient plant recorded a value 18.7%. The energy penalties ($f_1$) to capture and store the CO$_2$ from these two plants with a modern temperature-swing separation system are 34% and 52%, respectively. From the distribution of thermal efficiencies, $f_1$ and the additional fuel requirement ($f_2$) associated with converting all or some of the U.S. coal-fleet to CCS can be calculated.

Fig. 7(b) shows the distribution of $f_2$ for the U.S. installed base with 20% available-waste-heat recovery. That distribution yields a spread in $E_{\text{bc}}$ between 78 and 96 kJ/mol because less efficient plants have more available waste heat. Fig. 7(b) shows the corresponding $f_2$ distribution, which spreads from 0.57 to 1.01 with a mean value of 0.66 and standard deviation of 0.05. Converting the entire PC installed base to CCS while keeping its electrical work output constant would require an additional $\sim 460$ million tonnes of coal annually (assuming an average energy content of 25 GJ/(tonne coal)). Alternatively, the energy penalty could be manifest in a decreased plant output. In that case, the power output of the U.S. coal fleet would drop by $\sim 78$ GW.

Under the 20% available-waste-heat recovery assumption, the difference between retrofitting the most efficient plants for CCS and retrofitting the least efficient plants is significant [Fig. 7(b)]. If the 10 most efficient plants were retrofitted to capture and store 80% of their CO$_2$, then an additional 6.5 million tonnes of coal would be required and 27 million tonnes of CO$_2$ emissions would be eliminated annually. Thus, the CO$_2$ abatement effectiveness (i.e., the mass ratio of CO$_2$ eliminated to additional coal required) for the top 10 plants is 4.1. On the other hand, retrofitting the 10 least efficient plants would require 2.4 million tonnes of additional coal and would only eliminate 6.8 million tonnes of CO$_2$ annually yielding a CO$_2$ abatement effectiveness of 2.8. Thus, retrofitting the 10 most efficient PC plants for CCS would eliminate 46% more CO$_2$ emissions per unit of additional coal than retrofitting the 10 least efficient plants.

These calculations were repeated for the most and least efficient 10% and 25% of current PC plants (Table 2). The columns in Table 2 reveal how the CO$_2$ abatement effectiveness depends on the fraction of available-waste heat that is recovered from retrofitting a particular ensemble of the most efficient plants versus the equivalent ensemble of least efficient plants.

### Fig. 7

(A) The thermal efficiency distribution of an ensemble of 420 large U.S. coal-fired power plants. These plants produced the equivalent of 218 GW of constant electric power in 2007 constituting 96% of all U.S. coal-fired power output. (B) The distribution of additional fuel requirements ($f_2$) is calculated from the power-plant efficiency distribution. From this distribution, the total additional fuel required is calculated to be 530 million tonnes of coal. These calculations assume, $\eta_{\text{sep}} = 40\%$, $\eta_{\text{w}} = 20\%$, $\eta_{\text{com}} = 65\%$, and the national average coal heat content (25 GJ/(tonne)).
The CO2 abatement effectiveness (i.e., the mass ratio of CO2 eliminated to additional coal required) for 6 ensembles of U.S. coal plants: The ensembles are organized by each plant’s reported thermal efficiency. The first row labeled ‘Top 2%’ is for the most efficient 2% of U.S. coal plants. Each column assumes a different value for \( \eta_{w2nd} \), which is the fraction of available-waste heat that is harnessed for separation. In the \( \eta_{w2nd} = 0 \% \) case, retrofitting the most efficient 10% of plants will eliminate nearly 30% more CO2 per unit of additional coal than retrofitting the least 10% of plants. As the \( \eta_{w2nd} \) increases, then the gap in CO2 abatement effectiveness decreases because less efficient plants have a greater amount of available waste heat.

<table>
<thead>
<tr>
<th>CO2 abatement effectiveness ( (\eta_{w2nd} = 0%) )</th>
<th>CO2 abatement effectiveness ( (\eta_{w2nd} = 20%) )</th>
<th>CO2 abatement effectiveness ( (\eta_{w2nd} = 40%) )</th>
</tr>
</thead>
<tbody>
<tr>
<td>Top 2% (top 10)</td>
<td>3.3</td>
<td>4.1</td>
</tr>
<tr>
<td>Top 10%</td>
<td>3.1</td>
<td>4.0</td>
</tr>
<tr>
<td>Top 25%</td>
<td>3.0</td>
<td>3.9</td>
</tr>
<tr>
<td>Bottom 25%</td>
<td>2.6</td>
<td>3.4</td>
</tr>
<tr>
<td>Bottom 10%</td>
<td>2.4</td>
<td>3.1</td>
</tr>
<tr>
<td>Bottom 2% (bottom 10)</td>
<td>2.1</td>
<td>2.8</td>
</tr>
</tbody>
</table>

The dependence of the energy penalty and the CO2 abatement effectiveness on base-line efficiency derives from the primary energy required for compression and from the available waste heat. More efficient power-plants have lower \( E_{mc} \) values, but they also have less available waste heat. Fig. 6(b) reveals the sensitivity of \( f_2 \) to both \( \eta_{pp} \) and \( \eta_{w2nd} \). That figure indicates that once \( \eta_{w2nd} \) exceeds 50\%, \( f_2 \) is essentially independent of power-plant efficiency. On the other hand, if \( \eta_{w2nd} = 0 \% \), then both Table 2 (column 4) and Fig. 6(b) reveal that retrofitting the most efficient plants is a substantially more effective method of CO2 emission abatement. Table 2 reveals a narrowing of the CO2 abatement effectiveness with increasing available-waste-heat recovery, but even in the high \( \eta_{w} \) scenario, retrofiting the most efficient PC plants is nevertheless measurably more effective than retrofitting the least efficient plants.

The financial costs of CCS are tightly related to the energy penalty (Fig. 1). Indeed, it is worth noting in Fig. 1 that while new PC construction appear superior to PC retrofits when measured by the common metric of dollars per tonne of CO2 avoided (Fig. 1(a)), retrofits and new construction are about equal when measured by the more relevant metric of cost of electricity from a CCS power plant (Fig. 1(b)). That is partially the result of lower fixed costs associated with plants that have been fully or partially amortized.

The correlation between the costs of CCS and the energy penalty coupled with the variance in expected energy penalties (Fig. 7(b)) suggests an optimal CCS deployment strategy. The cheapest path to drastically reduce CO2 emissions from electricity production will combine the selective retrofitting of the most efficient PC plants with the closing of the least efficient plants. Overall, our analysis strongly suggests that the supply curve for retrofitting PC plants for CCS is a function of the power-plant’s baseline thermal efficiency.

It should be noted, however, that the relationship between base-line efficiency and the energy penalty assumes that the power plant itself is providing the compression work. Other configurations are possible. For example, a dedicated natural-gas-fired compressor or even a wind turbine could provide the necessary compression work. In those scenarios, \( f_1 \) would manifest as the reduced power output from either natural gas or wind; \( f_2 \), however, would be lower because the CO2 intensity of both gas and wind are lower than that of coal.

To demonstrate the potential value of available-waste-heat recovery, we calculate \( f_2 \) for the entire U.S. coal fleet with and without 1/3 available-waste-heat recovery. Retrofitting the entire U.S. coal fleet with zero available-waste-heat recovery would require additional ~600 million tonnes of coal annually. If, on the other hand, 1/3 of the available waste heat were productively used for separation, then the additional fuel requirement would drop from ~600 million to ~390 million tonnes of coal annually. Alternatively, if the energy penalty were manifest in a reduced power output, then with zero available-waste-heat recovery an additional ~92 GW of CO2-free base-load power would be required to make up for the decrease in power output. With 1/3 available-waste-heat recovery, however, the additional power requirement would drop from ~91 GW to ~69 GW.

Improving end-use electrical efficiency is an additional path through which the CCS energy penalty could be offset. This path is intriguing because the total U.S. smoothed power output was ~472 GW in 2007, indicating that increasing end-use electrical efficiency of between 15% and 20% would be sufficient to make up for the decrease in power output after retrofitting the installed PC base for CCS. That would yield a ~65% reduction in CO2 emissions from the power sector while not requiring any additional power-generation capacity to be build or any additional coal to be burned. The remaining 35% would come primarily from natural gas as well as a little from coal as we assumed 80% CO2 capture. This approach may be feasible as California has been able to keep its per capita electricity use constant for the past 30 years, while average per capita electricity use in U.S. grew by nearly 50%.

Conclusion

Achieving substantial reductions in CO2 emissions requires either shutting down a large fraction of the current installed base of coal-fired power plants or retrofitting those plants for CCS. Previous studies have estimated that the additional fuel required (\( f_2 \)) to maintain constant work output for a PC retrofit is between ~50% and 80%. An analysis of the thermodynamic limit indicates those values might be improved by harnessing more of the available waste heat and by improving the 2nd-law efficiency of temperature-swing separation systems. It appears difficult, however, to improve \( f_2 \) for post-combustion capture to below ~25% in practice. Our most likely efficiency scenario indicates that offsetting the energy penalty incurred from capturing and storing 80% of the U.S. coal fleet’s CO2 emissions will require either an additional ~390-600 million tonnes of coal, an additional ~69–92 gigawatts of CO2-free-base-load power, or a 15%–20% reduction in overall electricity use.
Nomenclature list with some characteristic values

<table>
<thead>
<tr>
<th>Symbol</th>
<th>Description</th>
</tr>
</thead>
<tbody>
<tr>
<td>( n_{CI} )</td>
<td>The number of moles of CO₂ in state ( i ) (state 1: ( \sim 0.12 ) moles CO₂ per mole of flue gas)</td>
</tr>
<tr>
<td>( n_{NI} )</td>
<td>The number of moles of N₂ in state ( i ) (state 1: ( \sim 0.80 ) moles N₂ per mole of flue gas)</td>
</tr>
<tr>
<td>( P_i )</td>
<td>The pressure of state ( i ) [Pa] (( \sim 10^{10} ) Pa at the surface, ( \sim 10^7 ) Pa in the reservoir)</td>
</tr>
<tr>
<td>( g )</td>
<td>The gravitational acceleration [m/s²]</td>
</tr>
<tr>
<td>( L_d )</td>
<td>The depth of CO₂ injection [( \sim 1000 ) m]</td>
</tr>
<tr>
<td>( L )</td>
<td>Length of pipeline</td>
</tr>
<tr>
<td>( \rho_w )</td>
<td>The density of H₂O [( \sim \text{kg/m}^3 )] (( \sim 1000 ) kg/m³)</td>
</tr>
<tr>
<td>( \rho_Ci )</td>
<td>The density of CO₂ in state ( i ) [( \text{kg/m}^3 )] (( \sim 2 ) kg/m³ at the surface, ( \sim 400-600 ) kg/(m³) in the reservoir)</td>
</tr>
<tr>
<td>( \rho_Ni )</td>
<td>The density of N₂ in state ( i ) [( \text{kg/m}^3 )] (( \sim 1.2 ) kg/m³)</td>
</tr>
<tr>
<td>( h_{CI} )</td>
<td>The molar enthalpy of CO₂ in state ( i ) [kJ/mol]</td>
</tr>
<tr>
<td>( h_{NI} )</td>
<td>The molar enthalpy of N₂ in state ( i ) [kJ/mol]</td>
</tr>
<tr>
<td>( s_{CI} )</td>
<td>The molar entropy of CO₂ in state ( i ) [kJ/(K mol)]</td>
</tr>
<tr>
<td>( s_{NI} )</td>
<td>The molar entropy of N₂ in state ( i ) [kJ/(K mol)]</td>
</tr>
<tr>
<td>( m_C )</td>
<td>The molar mass of CO₂ [kg/mol]</td>
</tr>
<tr>
<td>( m_N )</td>
<td>The molar mass of N₂ [kg/mol]</td>
</tr>
<tr>
<td>( T_a )</td>
<td>Temperature of the MEA absorber unit [( \sim 320-350 ) K]</td>
</tr>
<tr>
<td>( T_S )</td>
<td>Temperature of the MEA stripper unit [( \sim 400 ) K]</td>
</tr>
<tr>
<td>( T_L )</td>
<td>Temperature of the environment [( \sim 293 ) K]</td>
</tr>
<tr>
<td>( T_H )</td>
<td>Temperature of the steam working fluid [( \sim 1000 ) K]</td>
</tr>
<tr>
<td>( f_1 )</td>
<td>The energy penalty</td>
</tr>
<tr>
<td>( f_2 )</td>
<td>The fraction of additional fuel required to maintain the constant power output</td>
</tr>
<tr>
<td>( V )</td>
<td>Total swept out pore volume [m³]</td>
</tr>
<tr>
<td>( v_i )</td>
<td>Molar volume of state ( i ) [m³/mol]</td>
</tr>
<tr>
<td>( G )</td>
<td>The Gibbs free energy [kJ]</td>
</tr>
</tbody>
</table>

Primary Energy:

\( E_a \) The primary energy required to separate the CO₂ from the flue gas [kJ] |
\( E_b \) The primary energy required to compress the concentrated CO₂ to reservoir pressure [kJ] |
\( E_c \) The primary energy required to separate CO₂ from the water (thermodynamic limit <1 kJ/mol) |
\( E_{sep} \) The incremental primary energy required to separate CO₂ from the flue [kJ] |
\( E_S \) The total primary energy required for sequestration [kJ] |
\( E_w \) The quantity of waste heat that can—in principle—be used in separation [kJ] |
\( E'_w \) The total waste heat produced [kJ] |
\( E_L \) The minimum quantity of heat transferred to the environment [kJ] |
\( E'_L \) The actual quantity of heat transferred to the environment [kJ] |
\( E_H \) The primary energy content of the fuel [kJ] |

Efficiencies

\( \eta_{leak} \) Ideal separation efficiency (\( \eta_{leak} = \eta_{com} \) for pressure swing, \( \sim 30\% \) for temperature swing) |
\( \eta_{sep} \) 2nd-law separation efficiency (\( \sim 50\% \)) |
\( \eta_c \) The power-plant Carnot efficiency (\( \sim 70\% \)) |
\( \eta_{hp} \) The power plant efficiency (\( 25\% \)-\( 45\% \)) |
\( \eta_{com} \) Isothermal compression efficiency (\( 65\% \)) |

Other Parameters

\( X_{Ni} \) The mole fraction of N₂ in state \( i \) (state 1: \( \sim 80\% \)) |
\( X_{CI} \) The mole fraction of CO₂ in state \( i \) (state 1: \( \sim 12\% \)) |

References


29. B. A. Oyenekan, Modeling of Stripper for CO₂ Capture by Aqueous Amines, Department of Chemical Engineering, The University of Texas at Austin, Austin, Texas, 2007.


