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Coherence and Raman Sideband Cooling of a Single Atom in an Optical Tweezer

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We investigate quantum control of a single atom in a tightly focused optical tweezer trap. We show that inevitable spatially varying polarization gives rise to significant internal-state decoherence but that this effect can be mitigated by an appropriately chosen magnetic bias field. This enables Raman sideband cooling of a single atom close to its three-dimensional ground state (vibrational quantum numbers $n_r = 0.01$, $n_z = 8$) even for a trap beam waist as small as $w = 900$ nm. The small atomic wave packet with $\delta x = \delta y = 24$ nm and $\delta z = 270$ nm represents a promising starting point for future hybrid quantum systems where atoms are placed in close proximity to surfaces.

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Single atoms in “optical tweezer” traps [1] are a promising resource for various applications in quantum science and engineering. They can be individually moved [2], manipulated [3,4], read out [5,6], and used to implement quantum gates [7,8], in a manner similar to trapped ions. At the same time, they may be strongly coupled to photonic [9,10], plasmonic [11], or other solid-state systems [12–14], opening a new frontier for the realization of quantum networks and hybrid quantum systems. These intriguing applications require trapping single ultracold atoms near surfaces at distances well below an optical wavelength. While this is challenging for ions [15] and magnetically trapped atoms [12,16], it is achievable with neutral atoms in optical dipole traps.

An optical tweezer can be efficiently loaded with a single atom from a magneto-optical trap by making use of a light-induced two-body loss process (collisional blockade) [1,17]. The temperature of an atom loaded in this way is in the range of 30–200 $\mu$K [3,4,7,8,10,18,19], at which point the atom has a characteristic root-mean-square (rms) spatial extent of $\delta r \approx 200$ nm and $\delta z = 1$ $\mu$m in the radial and axial directions, respectively. This spatial spread is an impediment in several current experiments [7,20], while the elevated temperature reduces the coherence time [3,4,7,8,19]. Moreover, interfacing the atom to the near field of a solid-state structure requires much stronger confinement, as in this case the atom must be localized on the scale of a reduced resonance wavelength $\lambda/2\pi = 120$ nm (for Rb).

These applications require significant improvements in laser cooling and coherent manipulation. Raman sideband cooling is a powerful technique to control atomic motion, as was demonstrated previously for ions and atomic ensembles in larger traps [21–25]. Coherent Raman transitions between two stable ground states that change the atom’s vibrational level can be used to remove energy, combined with an optical pumping process to remove entropy and complete the cooling cycle [Fig. 1(a)].

However, in very tightly confining traps with beam waist $w \sim \lambda$, polarization effects associated with the breakdown of the paraxial approximation can strongly impede coherent manipulation and cooling. Such tightly focused beams exhibit a longitudinal polarization component, which even for linearly polarized input fields results in spatially varying elliptic polarization [26–29]. The corresponding atomic-state-dependent trapping potentials reduce atomic coherence, induce force fluctuations, and impair cooling [30]. These effects are present not only in optical tweezers but also at subwavelength distances from dielectric boundaries [10,31] and in projected optical lattices to be used for many-body quantum simulation [32].

In this Letter, we present a detailed study of the longitudinal polarization component of a dipole trap formed by a high-numerical-aperture lens, demonstrate how the associated effect on a trapped atom can be partially compensated using a properly oriented magnetic bias field, and apply these results to perform Raman sideband cooling of a single atom. After cooling, the atom is in the ground state along the two radial directions ($\tilde{n}_r = 0.01^{+0.06}_{-0.01}$) and occupies just a few quantum states ($\tilde{n}_a = 8.1(8)$) in the axial trap direction. The corresponding rms size of the atomic wave packet is given by the ground-state length of $\delta r = 24$ nm in the radial directions and a thermal extent $\delta z = 270$ nm in the axial direction. This represents a 100-fold reduction in spatial volume and a reduction by $10^4$ in phase-space volume over the starting conditions.

The longitudinal polarization component can be understood in the framework of ray optics [see Fig. 1(b)]. Light entering a lens consists of parallel rays with transverse linear polarization. Upon refraction, the polarization of each ray must also deflect to remain transverse to the ray [26]. In the diffraction-limited volume around the focus, all rays interfere and the resulting field is elliptically polarized. Following Fig. 1(b), two features emerge near the focus: The polarization vector is rotating in the plane set by
FIG. 1 (color). (a) Relevant levels and transitions in $^{87}$Rb. The eigenstates of the harmonic potential for the ground state are indicated with dashed lines. Atomic levels are defined in the $|F, m_F\rangle$ basis. See the text for beam orientations and polarizations. (b) The origin of the elliptical polarization near the focus (see the text). In the figure, (C)CW refers to (counter-) clockwise rotation of the polarization vector. (c) A cut through the focal plane for gradient. The fit is to the model described in the text: $dC_x=dy$. (d) Dephasing rate between the states $|1\rangle$ and $|2\rangle$ as a function of bias field, with $\lambda_p = 815$ nm. The improvement at large bias fields is due to suppression of the polarization gradient. The fit is to the model described in the text: $\eta_0+\eta_{\text{circ}}$. The background shading shows the Gaussian intensity profile for comparison. (d) Dephasing rate between the states $|1\rangle$ and $|2\rangle$ as a function of bias field, with $\lambda_p = 815$ nm. The improvement at large bias fields is due to suppression of the polarization gradient. The fit is to the model described in the text: $\eta_0+\eta_{\text{circ}}$. The background shading shows the Gaussian intensity profile for comparison.
is extracted from a Ramsey-type measurement, using a fit to the function introduced in Ref. [30].

At two different trap wavelengths $\lambda_T$, we fit $1/T_2^\text{rad} = 1.03(\eta_0 + \eta_{\text{circ}} + \eta_{\text{rec}})(\hbar g/k_B T/2h)$. The only free parameters are the degree of circular polarization in the incident dipole trap beam due to uncompensated birefringence along the beam path ($\eta_{\text{circ}}$) and the strength of the effective field gradient $B'_r$. The temperature is determined independently ($T = 40 \mu\text{K}$ for this measurement; see below for the technique). $\eta_0$ reflects the different trapping potentials for $F = 1$ and $F = 2$ atoms due to the finite trap detuning.

At $\lambda_T = (802.815) \text{nm}$, we find $B'_r = (2.4, 1.4) \text{G}/\mu\text{m}$, and thus $dC_x/dy = [0.46(6), 0.54(3)]/\mu\text{m}$, in reasonable agreement with our estimate of 0.57 $\mu\text{m}$.

Having developed a detailed understanding of trap-induced decoherence in this system, we now turn to Raman sideband cooling. We use three orthogonal running-wave fields to drive Raman transitions, labeled $R1-R3$ [Fig. 1(a)]. $R1$ propagates antiparallel to the dipole trap (−$\hat{z}$) and is circularly polarized to drive $\sigma_-$ transitions. $R2$ propagates along $\hat{x}$ and is circularly polarized; $R3$ propagates along $\hat{y}$ and is linearly polarized along $\hat{x}$. Optical pumping to the $|2\rangle$ state is provided by circularly polarized beams copropagating with $R1$, addressing the $F = 1 \rightarrow F' = 2$ and $F = 2 \rightarrow F' = 2$ transitions on the $D2$ line. The frequencies of the lasers are set to the measured resonances in the dipole trap, which are shifted by $\sim 30 \text{MHz}$ from the resonances in free space; the intensities are about 100 times less than saturation. We measure the $F = 1$ population by pushing out any atom in $F = 2$ using a circularly polarized beam along the optical pumping path that is resonant with the $F = 2 \rightarrow F' = 3$ transition on the $D2$ line, then measuring whether the atom has remained trapped by turning the molasses back on.

In a typical experiment, we load a single atom from the magneto-optical trap into the dipole trap with a depth of 1.6 mK at zero bias field, then decrease the trap depth to 0.82 mK while ramping the bias field $B_z$ up to 7.5 G. Lowering the trap depth serves to increase the coherence time while leaving the trap frequencies high enough that sideband cooling is still achievable, with $(\omega_r, \omega_a) = 2\pi \times (100, 15.6) \text{kHz}$. All temperatures reported in this Letter are measured in the 0.82 mK deep trap. We cool the atoms in the following sequence: We first apply the $R2$ and $R3$ beams (Fig. 1) and the optical pumping beams together for 10 ms to continuously cool the radial modes; then, we perform ten cycles consisting of 2 ms of axial cooling using the $R1$ and $R2$ beams, followed by 4 ms of radial cooling using the $R2$ and $R3$ beams again. This sequence prevents the radial modes from heating while the axial cooling proceeds.

The parameters for the first radial cooling phase are optimized by measuring the temperature using a release-and-recapture technique [37]. These data, shown in Fig. 2(a), are fit using a Monte Carlo simulation [18]. The initial kinetic energy per dimension $K$ is such that $2K/k_B = 52 \mu\text{K}$. The measurement after cooling yields anisotropic kinetic energies of $2K_r/k_B = 2.4(1) \mu\text{K}$ in the radial direction and $2K_a/k_B = 158(14) \mu\text{K}$ in the axial direction (the release-and-recapture technique is only weakly sensitive to the axial mode). The fitted kinetic energies represent the global minimum in $\chi^2$ over the entire space of three independent energies for each axis, including unphysical temperatures less than the ground-state energy $\hbar \omega/2k_B = 2.4 \mu\text{K}$ for the radial modes.

The agreement of the measured kinetic energy with that of the zero-point motion suggests that we have reached the radial ground state after this cooling phase alone. The radial cooling works best with a two-photon Rabi frequency $\Omega_{R2,R3} = 2\pi \times 17 \text{kHz}$ and a detuning of $-\omega_r = -2\pi \times 100 \text{kHz}$ from the two-photon resonance.

To characterize the axial temperature independently after the radial cooling, we measure the Doppler width of the $|2\rangle$ to $|1\rangle$ transition when driven with the $R1$ and $R2$ beams. The wave vector $\Delta k_{12} = k_{R1} - k_{R2}$ has a projection onto the axial and radial directions, but the Doppler profile should mostly be sensitive to the axial mode here since the radial degrees of freedom are already cold. After the first stage of radial cooling, we measure a kinetic energy of $2K_a/k_B = 129(19) \mu\text{K}$ [Fig. 2(b)]. After implementing and optimizing the axial cooling, we

FIG. 2. (a) Release-and-recapture temperature measurement. The closed and open circles show measurements before and after radial cooling, respectively. A Monte Carlo model yields kinetic energies $K$ such that $2K/k_B = 52(4) \mu\text{K}$ before cooling, and $(2K_r/k_B, 2K_a/k_B) = [2.4(1), 158(14)] \mu\text{K}$ after cooling.

(b) Doppler measurement of the axial kinetic energy before and after cooling the axial mode. (b) After radial cooling only, $2K_a/k_B = 129(19) \mu\text{K}$. (c) After radial and axial cooling, $2K_a/k_B = 8.1(1) \mu\text{K}$. 133001-3
We also resolve the axial motional sidebands using the \( R_1 \) and \( R_2 \) beams at very low power and observe a spectrum with nine peaks that is slightly asymmetric [Fig. 3(b)]. We find that the ratios of the measured peak heights correspond very well to a thermal distribution \( \rho_{nn} \propto \exp(-n/\tilde{n}_a) \) with a mean vibrational number \( \tilde{n}_a = 8.1(8) \). The corresponding energy \( (\tilde{n}_a + 1/2)\hbar\omega_a = 6.5 \mu\text{K} \times k_B \) is similar to the result of the Doppler measurement above.

Several properties of the cooled atom are worth noting. The heating rate for the radial degrees of freedom is very low, less than \( \Delta \tilde{n} < 0.3 \) over 200 ms. Also, we observe no heating while translating the atom over distances \( \sim 20 \mu\text{m} \) in \( \sim 10 \text{ms} \) using a scanning galvanometer mirror. Decreasing the Rabi frequency \( \Omega_{R_1,R_2} \) and detuning during the last cooling phase does not decrease the final axial temperature. This is possibly due to the fact that we cannot separately address the axial mode or to our choice to optically pump along the axial direction, resulting in more heating along that direction. We are not aware of any fundamental effects that would prevent cooling to the ground state in this system.

It may be possible to extend the demonstrated method to perform high-fidelity state detection [5,6] while cooling within one hyperfine state and collecting optical pumping photons. Furthermore, it should also be possible to cool small ensembles of atoms held in arrays of traps [17] or together in a single trap. In the latter case, Raman cooling is advantageous compared to an optical molasses, in that the detuning of the optical pumping beam can be chosen over a wide range, allowing the effects of light assisted collisions [38] and heating due to rescattered photons [24] to be reduced.

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Note added.—After completion of this work, we have become aware of a related demonstration of Raman sideband cooling in an optical tweezer [39].

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