The long evolution of vascular plants has resulted in a tremendous variety of natural networks responsible for the evaporatively driven transport of water. Nevertheless, little is known about the physical principles that constrain vascular architecture. Inspired by plant leaves, we used microfluidic devices consisting of simple parallel channel networks in a polymeric material layer, permeable to water, to study the mechanisms of and the limits to evaporation-driven flow. We show that the flow rate through our biomimetic leaves increases linearly with channel density \( (1/d) \) until the distance between channels \( (d) \) is comparable with the thickness of the polymer layer \( (\delta) \), above which the flow rate saturates. A comparison with the plant vascular networks shows that the same optimization criterion can be used to describe the placement of veins in leaves. These scaling relations for evaporatively driven flow through simple networks reveal basic design principles for the engineering of evaporation–permeation-driven devices, and highlight the role of physical constraints on the biological design of leaves.

Leaves are the main pumps for the upward movement of fluids in plants, which takes place in a network of conduits collectively called xylem (1). On warm, sunny days, the leaves of a large tree can transport hundreds of liters of water from the soil, corresponding to several milliliters per day per leaf. Although photosynthesis, rather than water loss, is the desideratum of leaves, the shared diffusional pathway for \( \text{CO}_2 \) and water vapor and the regulation of this pathway by water status in the leaf implies that these two fluxes are closely coupled. Indeed, maximum rates of photosynthesis scale linearly with the transpiration rate driven by diffusion are inversely related to the distance over which the gradients operate. This analogy allows us to mimic water transport in leaves using a simple physical model in which water moves from supply channels (veins) to the leaf surface solely by diffusion.

To examine the design constraints on evaporatively driven hydraulic systems (Fig. 1 and Experimental Methods) we use microfluidic systems that use permeation-driven flows (10, 11). The ability to manufacture and manipulate complex microfluidic geometries (12) using poly(dimethylsiloxane) (PDMS) allowed us to explore analogies to leaves in a quantitative, measurable way using our knowledge of water permeation through PDMS in microchannels. As shown recently (11), the transport mechanism is dominated by molecular diffusion of vapor driven by concentration gradients. Water molecules move in small clusters through the hydrophobic PDMS network from one pore to another over a scale of a few angstroms, just as water vapor diffuses in air, but with a much smaller diffusion coefficient because of the presence of hydrophobic barriers. This simple picture is valid in the immediate vicinity of a single channel of water in the PDMS leaf analog.

For the case of multiple sources or channels shown in Fig. 1, if the thickness of the PDMS layer is small compared with the spacing between the channels, i.e., \( \delta \ll d \), each channel provides an evaporative surface flux that is independent of its neighbor. However, when the layer thickness is large compared with the distance between channels, \( \delta \gg d \), the channels collectively act as a single one. With these simple ideas in mind, our goal was to determine, using both experiment and theory, the scaling relations that constrain evaporatively driven transport through simple devices.
Results

Effects of Environment and Geometry. By varying the external humidity imposed at the PDMS surface, we find that our devices behaved as free evaporative surfaces, with flow rates proportional to the relative dryness (= 1 – RH) imposed at the surface (Fig. 2). However, the rate of water movement was sensitive to channel density; at low relative humidities, a greater density of channels resulted in higher flow rates. The minimum channel density that can fully satisfy the evaporative demand is the key design parameter for natural and artificial pervaporative pumps.

An oversupply of channels represents a significant inefficiency, both in terms of construction costs and the consumption of space that might otherwise be devoted to other functions. At low channel density (δ ≪ d, where d is the distance between channels), each channel functions independently such that the total flow rate is proportional to channel number. The concentration of water within the PDMS layer decreases toward the air, with the exact concentration profile dependent on channel shape and size. Except in the region very close to the channel walls, the concentration decreases logarithmically as a function of distance from the channel, as expected from the solution of the Laplace equation (see Model: Analytical Approximations and Model: Numerical Simulation). The total flow rate (J_{T}) is the sum of the flux from each channel, which can be written as

\[ J_{T} \approx \sum_{i} \frac{\pi D (c_{0} - c_{1}) L}{\ln \left( \frac{\delta + h}{a} \right)} \frac{D (c_{0} - c_{1})}{d} L W, \delta \ll d, \]

where D is the diffusivity of water vapor in PDMS, c_{0} and c_{1} are, respectively, the concentrations at the channel and evaporative surface, L is the channel length, h is the channel height, d is the distance between channels, W is the width of the microfluidic device, a is the equivalent radius of a semicircular channel (see Model: Analytical Approximations and Model: Numerical Simulation), and N = Wd is the number of channels.

Conversely, for high channel density (δ ≫ d), the vapor concentration decreases linearly toward the PDMS–air interface such that the flow rate is independent of channel density and the flux is then

\[ J_{T} = \frac{D (c_{0} - c_{1})}{\delta} L W, \delta \gg d. \]

In Fig. 3, we show that flow rate through artificial leaves of varying thickness is in good agreement with theory; for a given thickness, it increases with channel densities and eventually saturates. Furthermore, the maximum flow rate J_{n} ∝ δ^{-0.92} (Fig. 4), close to the simple prediction for diffusively limited transport in a 1D medium.

Optimization. Our preceding results suggest that for a given PDMS layer thickness, there is a characteristic channel spacing (d_{c}) below which there is no additional gain in the flow rate because the system behaves as if it were effectively a quasi one-dimensional medium. Equating the flux when d ≫ δ (Eq. 1) to the flux when d ≪ δ (Eq. 2), with n = Wd, we find

\[ d_{c} = \frac{\pi \delta}{\ln \left( \frac{\delta + h}{a} \right)}. \]

In Fig. 3, one can check that the characteristic channel spacing above which the flow saturates for each thickness is in good agreement with Eq. 3. In the general case, corresponding to different channel shapes and dimensions, this exact result is replaced by a simple, intuitive scaling law,

\[ d_{c} \sim \delta, \]

Additional mathematical and experimental support is provided in the supplementary material.
as expected on dimensional grounds for the transition from two-dimensional behavior to one-dimensional behavior that occurs when the channel spacing becomes comparable with the PDMS layer height and the channels act collectively rather than as individuals. Indeed, not only is there no gain in flow by adding more channels for pervaporative pumps of a given thickness, but when the channels are too close to each other, the surface area over which PDMS adheres with glass decreases so that the maximum pressure sustainable by the device is lowered. This suggests that for this particular artificial leaf, a compromise between adhesive strength and flow rate determines the optimal channel spacing. More generally, if there is a cost associated with making veins (e.g., ensuring that adhesion is large enough, or in a biological context, a metabolic cost), the scaling relation Eq. 4 provides a simple optimization principle for the design of pervaporative devices.

Real Leaves. The diversity of leaf venation found in nature provides an ideal opportunity to test the generality of the scaling relations demonstrated for artificial leaves. Most leaves exhibit a hierarchical and reticulate venation, with the larger veins serving as major supply lines for the smaller veins that serve as the principal sites where water moves from the xylem to the mesophyll. Because we are interested in the ability of the venation to supply the evaporative demand, rather than distribute water across the leaf as a whole, we focused on the placement of only the highest-order veins (7). Although there is a large literature on the shape (13) and sizes of venation patterns, there are few studies (7) that simultaneously report leaf thickness, and only a recent paper (14) provides information on the physically meaningful parameter of the distance from the xylem to the evaporative surface. We therefore measured both the average distance between the highest-order veins and the leaf evaporative surface ($d$) and the average distance between adjacent veins ($d'$), for 32 species characteristic of open, unshaded environments (see Experimental Methods). The fact that a leaf is reticulate or parallel has no bearing on either the determination of $d'$ or on the interpretation of these data using our model.

The venation of this diverse assemblage of leaves shown in Fig. 5 is well described by the simple law (Eq. 4), suggesting that leaves do not produce veins in excess of what are needed to allow maximum rates of water delivery to the leaf surface. Factors contributing to this optimum include both the costs of constructing the veins and the utilization of space that could otherwise be filled with photosynthetic cells. A simple fit gives $d = 1.08 \delta$, $R = 0.989$, quantitatively consistent with Eq. 4. However, one has to keep in mind that our model only provides a scaling relationship, and a calculation of the exact numerical prefactor of $\delta$ remains beyond the ability of any simple model unless we understand the multiple parameters that affect the design of real leaves. Nevertheless, our scaling law provides a theoretical framework for the recently reported inverse correlation between photosynthesis and distance for water movement between veins and stomata (14).

Discussion

In this study, we have explored only the simplest topology of channels in our artificial leaves, ignoring variations in the morphology and architecture that characterize the complex venation systems of real leaves. The general scaling principles that our biomimetic leaves follow in response to the constraints from the evaporative surface reflects the design requirements for optimal channel placement. Taken together with the strong correlation between density of veins and their distance to the leaf surface in real leaves, this suggests that design features oriented toward optimal utilization of veins in providing water for evaporation are universal and thus independent of the morphological diversity among leaves. Because of the analogy of transpiration rate variation with distance due to active stomatal control and of diffusion in PDMS leaves, our study also provides further evidence of the homeostatic regulation of leaf water status via stomatal control. For natural systems, the existence of architectural optima describing the placement of veins in leaves provides an important foundation for understanding leaf function (5) as well as the physical/physiological limits imposed by the distribution system itself on development (8). For biomimetic systems, our study points toward principles for their design based on an understanding of the underlying physical limits.

Methods and Models

Experimental Methods. Microchannels were fabricated by using PDMS, as described in ref. 15, with 10:1 PDMS:cross-linker mixture (Sylgard 184), spin-coated onto a negative photoresist SU-8 mold, degassed for 15 min at a 4-inch Hg vacuum pressure, and cured at 65°C for 4 h. The channels (width $w = 15 \mu m$, height $h = 30 \mu m$) were made in a thin layer bonded to a glass slide by plasma treatment, its thickness adjusted by the spin-coating speed. A thick layer was

In this reference, qualitative correlation is made between vein density and tissue proportions, which is not necessarily related to $\delta$. 

![Fig. 4. Saturation flow rate (asymptotic limit for very small vein spacing) at 50% RH as a function of thickness (log-log plot). The straight line is a power fit with an exponent of $-0.92$.](image)

![Fig. 5. Measured distance between veins ($) versus distance to evaporative surface ($) for different plant species. The solid line corresponds to the transition to the plateau regime $d_0 = \delta$.](image)
bonded by plasma treatment on top of the thin layer to make the inlet for the liquid and to act as a frame, avoiding wrinkles in the PDMS film. This layer presents a rectangular window the size of the evaporative surface. The channels were slowly filled with water, the air escaping through the PDMS, because the channels had sealed ends. The devices were stored in a wet environment to keep the channels filled with water because contact with air results in the PDMS surfaces regaining its hydrophobicity (16), which makes refilling them more difficult. The channels (length $L = 35 \text{ mm}$), were connected over a width $W = 30 \text{ mm}$ to a large inlet channel (100 $\mu\text{m}$ wide) that feeds them. The spacing between two channels $d \in [30 \mu\text{m} to 1,500 \mu\text{m}]$, and the PDMS thickness above the channel $\delta \in [45 \mu\text{m} to 760 \mu\text{m}]$ (Fig. 1). The channel density is then $n = 1/d$, and the number of channels over 30 mm is $n = W/d$.

The devices were connected below them (30 cm) to a beaker resting on the plate of a precision balance (Sartorius BP 211 D) to ensure that water was stable. Then, we may consider a 2D problem for the evaporation of water from a simple linear profile due to a constant gradient of concentration. We see that for $d \gg \delta$, this leads to a linear relation for the flux as a function of the number of channels:

$$J_x = \frac{n D_{c_0} - c_1}{L} \ln \left( \frac{\delta + h}{a} \right) W.$$  

Extrapolating Eq. 6 into the limit $d \ll \delta$ (or $N \rightarrow \infty$), we find that this expression is close to the one obtained from Eq. 2 (with a numerical coefficient in $[\ln(\delta + h)/a][\ln(\delta + h)/h]$). This is not surprising as the channels are no longer independent in this limit; nevertheless it still gives the appropriate form of a linear increase followed by a plateau as a function of channel density.

Model: Numerical Simulation. We used a finite element method implemented in Matlabc (pdeTool) to solve the steady state diffusion equation for the water vapor concentration in two dimensions: $\Delta c - 0$, using the particular geometry of our system. We assume that the concentration and fluxes are periodic in the $x$ direction. We further assume that there are no edge effects on the three sides of the unit cell of the artificial leaf, symmetry dictates that the transversal flux $J = -D(d/dx)$, while on the two sides of the unit cell of the artificial leaf, the vertical flux is assumed to be zero, $J = 0$. This lead to a linear relation for the flux as a function of the number of channels:

$$J_x = \frac{1}{\ln \left( \frac{\delta + h}{a} \right)} \frac{D_{c_0} - c_1}{(\delta + h)^2} W.$$  

with $c(x, z, t)$ being the concentration, $z$ being the coordinate along the direction of evaporation, and $x$ being the coordinate perpendicular to it, along the channel width. The associated boundary conditions are (i) at the air interface $c = c_1$, which we will implement eventually using the technique of images, and (ii) there is no flow on the boundaries $x = \pm d/2$ and $z = \pm d/2$.

For the general case of a periodic variation of the concentration in the PDMS, we can decompose the solution of the above equation in a Fourier series of the form $c(x, z) = \sum_{n} c_n \cos(n \pi x / a) \exp(-n^2 \pi^2 D / 4 t)$, where the first term describes variations due to the individual channels, and the second term is the simple linear profile due to a constant gradient of concentration. We see that for thicknesses larger than the distance between channel, $\delta > d$, the periodic terms vanishes and only the 1D term remains so that the concentration decreases linearly toward the air interface, showing why the transition be-

between a 2D-like and 1D-like diffusion occurs for $d \ll \delta$. When $d \gg \delta$, however, the channels are effectively decoupled, so that it is sufficient to consider a unit cell around one channel ($-d/2 < x < d/2$). Replacing the channel of rectangular cross-section with an effective channel with a semicircular one (for simplicity), of radius $a = (2 \sqrt{Wh}/4)$, yields the result that the concentration decreases logarithmically with distance from the channel. In reality, channels are rectangular, but their effective radius a can be determined by using numerical simulation (described later) and leads to very similar results for the range of values for $d/\delta$ we have used. Then, one finds that the flux at the film-air interface for a single channel is

$$J_x = \frac{D_{c_0} - c_1}{L} \ln \left( \frac{\delta + h}{a} \right) W.$$  