HNC in Protoplanetary Disks

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HNC IN PROTOPLANETARY DISKS

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ABSTRACT

The distributions and abundances of small organics in protoplanetary disks are potentially powerful probes of disk physics and chemistry. HNC is a common probe of dense interstellar regions and the target of this study. We use the Submillimeter Array (SMA) to observe HNC 3–2 towards the protoplanetary disks around the T Tauri star TW Hya and the Herbig Ae star HD 163296. HNC is detected toward both disks, constituting the first spatially resolved observations of HNC in disks. We also present SMA observations of HCN 3–2, and IRAM 30m observations of HCN and HNC 1–0 toward HD 163296. The disk-averaged HNC/HCN emission ratio is 0.1–0.2 toward both disks. Toward TW Hya, the HNC emission is confined to a ring. The varying HNC abundance in the TW Hya disk demonstrates that HNC chemistry is strongly linked to the disk physical structure. In particular, the inner rim of the HNC ring can be explained by efficient destruction of HNC at elevated temperatures, similar to what is observed in the ISM. To realize the full potential of HNC as a disk tracer requires, however, a combination of high SNR spatially resolved observations of HNC and HCN, and disk specific HNC chemical modeling.

Keywords: Astrochemistry – ISM: molecules – Protoplanetary Disks

1. INTRODUCTION

The disks around young stars contain the gas and dust from which planets form. From theory (Aikawa et al. 1996; Bergin et al. 2003) and observations (Rosenfeld et al. 2012; Qi et al. 2013), disks are characterized by steep temperature and radiation gradients both radially and vertically. These disk structures affect planet formation (Takeuchi et al. 1996; Kretke & Lin 2012), motivating the quest for new emission-line diagnostics that can trace physical structure gradients in disks.

This study focuses on HNC, a common high density gas tracer in the ISM. At low temperatures, HNC forms efficiently together with HCN through ion-molecule reactions. HNC and HCN are observed in a variety of different sources, but their relative abundances can vary by more than order of magnitude, which has been related to the physical environment where the emission originates. In pre-stellar and protostellar environments, the observed HNC/HCN abundance ratio is ∼1 at 10 K and decreases as temperature increases (e.g. Schilke et al. 1992; Pratap et al. 1997; Liszt & Lucas 2001; Jørgensen 2004; Padovani et al. 2011; Jin et al. 2013), suggesting that thermal reaction barriers play an important role in controlling this ratio. Work by Loison et al. (2014) has shown that the HNC/HCN ratio also depends on the elemental abundances of carbon at low temperatures, due to destructive reactions with atomic carbon. In less shielded regions, the HNC/HCN ratio also depends on the radiation present (Meijerink et al. 2004).

Nearby, protoplanetary disks, such as the one surrounding TW Hya, provide an interesting test bed for evaluating the relative impact of different environmental parameters on the HNC chemistry, due to their relatively well characterized geometries and density, temperature and radiation structures. The HNC distribution in disks is interesting in its own right. HNC and HCN are two of very few nitrogen-bearing molecules detected toward disks (Dutrey et al. 1997, 2007; Chapillon et al. 2012; Obers et al. 2013) and their distributions may therefore provide some of the best available constraints on the nitrogen reservoir. Cyanides and iso-cyanides are also important for some prebiotic chemistry scenarios (Powner et al. 2007). Observing HNC in disks has proven challenging, however. The intrinsic disk scales are small and the molecular line fluxes are weak compared to interstellar clouds and protostars. To date, HNC has only been detected in a single disk, DM Tau, (Dutrey et al. 1997) and no spatially resolved observations exist.

In this letter, we report the first interferometric detection of HNC in two disks, around the T Tauri star TW Hya (D = 54 pc Qi et al. 2004) and Herbig Ae star HD 163296 (D = 122 pc Qi et al. 2011), using the SMA. We combine these observations with new and previous SMA observations of HCN toward TW Hya and HD 163296, and new IRAM 30m observations of HNC and HCN 1–0 toward HD 163296, with the aim to constrain the HNC/HCN chemistry in disks.

2. OBSERVATIONS OF HNC AND HCN

The HNC 3–2 line at 271.981 GHz was observed toward the TW Hya disk (Δα = 11ʰ01ᵐ51.875ˢ, ∆δ = [dgraninger@cfa.harvard.edu]
peaks at 0.53 Jy km s$^{-1}$ beam$^{-1}$. The integrated HNC emission peaks are offset to the east and west of the continuum phase center, suggestive of a ring-like emission structure within the disk. In the HNC image the western flux peak is almost twice as strong compared to the eastern one, corresponding to $\sim 3\sigma$. Asymmetries of this magnitude may be artifacts of CLEANing process and based on visual inspection alone we are agnostic about its significance. We evaluate this aspect of the data when modeling the complex visibilities in §3.2.

To isolate the radial flux variation in the HNC emission, Fig. 4 (lower right panel) shows the deprojected azimuthally averaged HNC line flux, using an inclination angle of 6$^\circ$ and position angle of 155$^\circ$ (Qi et al. 2004; Williams & Cieza 2011). The resulting emission profile reaches a maximum at $\sim 0.5''$, indicating the presence of a small HNC hole at the center of the TW Hya disk. This agrees with the $\sim 1''$ separation between the two flux peaks in the HNC integrated emission image. The depth and size of the hole may both be affected by the image asymmetry, however. In §3.2 we evaluate the evidence for an HNC ring and quantify its location based on an analysis in the $(u, v)$-plane.

The HNC spectrum in Fig. 3 was extracted using a 2$''$$\times$ 2$''$ box. The integrated line flux is 1.2 Jy km s$^{-1}$ beam$^{-1}$ ($\sim 12\sigma$). Using the the integrated HNC flux from this study and the HCN flux from Qi et al. (2004), the disk-averaged HNC/HCN emission ratio is 0.18 $\pm$ 0.06, where the error is derived using the 1$\sigma$ rms for the two observations.

HNC is also detected toward HD 163296, but with a lower SNR. Figure 2 displays the 1.1 mm continuum, HNC, and HCN emission toward HD 163296 as well as the spectra of the molecules and the velocity structure of

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1 http://www.cfa.harvard.edu/~r-cqj/mircook.html
2 http://www.iram.fr/IRAMFR/GILDAS

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Figure 1. The 1.1 mm continuum and HNC 3–2 emission (upper panels) from TW Hya using the compact + extended configurations of the SMA. Continuum contours are [10,20,40,...] mJy and the contours for HNC are [4,6,8]σ with 1σ = 0.064 Jy km s$^{-1}$ beam$^{-1}$. The lower left panel displays the spectrum obtained from the HNC 3-2 image. The lower right panel displays the deprojected, azimuthally averaged flux as a function of radius.
HNC in Disks

Figure 2. 1.1 mm continuum (top left, contours=[20,40,80,...] mJy), integrated HNC 3–2 (top middle) and HCN 3–2 (top right) fluxes toward the disk around HD 162936 together with extracted spectra (bottom left) and integrated emission in two velocity bins (bottom middle and right). The HNC and HCN integrated flux contours are [2,3]σ (HNC) and [3,5,7,...]σ (HCN) with 1σ = 0.21 and 0.14 Jy km s\(^{-1}\) beam\(^{-1}\), respectively. The HNC and HCN contours in the two velocity bins are [3,4]σ (HNC) and [8,10,12,...]σ (HCN). The velocity integrated emission maps yields a maximum HNC and HCN flux of 0.71 Jy km s\(^{-1}\) beam\(^{-1}\) and 2.71 Jy km s\(^{-1}\) beam\(^{-1}\), respectively. The HNC emission toward the HD 163296 disk is suggestive of a ring viewed at relatively high inclination, but the low SNR inhibits any further spatial analysis.

Spectra were extracted using an elliptical mask, which increased the SNR compared to a square mask. The total integrated flux from these spectra is 0.66 Jy km/s beam\(^{-1}\) (\(\sim 4\sigma\)) and 4.47 Jy km/s beam\(^{-1}\) (\(\sim 9.5\sigma\)) for HNC and HCN, respectively. This yields an HNC/HCN ratio of 0.14 \(\pm 0.05\), similar to TW Hya.

The IRAM 30m observations of HNC and HCN 1–0 toward HD 163296 should be more sensitive to the coldest disk material (Fig. 3). The integrated intensities of HNC and HCN were found to be <0.033 K km s\(^{-1}\) (3σ upper limit) and 0.104 K km s\(^{-1}\), respectively. This yields an upper limit on the HNC/HCN ratio of <0.3, consistent with the SMA data.

3.2. HNC line modeling toward TW Hya

To constrain the spatial extent of the HNC emission in TW Hya, we follow the procedure outlined in Qi et al. (2008), Oberg et al. (2012, 2015) to extract molecular column density profiles. First, parametric HNC abundance models are defined with respect to the TW Hya disk temperature and density structures from Qi et al. (2013B). The line emission is modeled using the two-dimensional non-LTE Monte Carlo radiative transfer code RATRAN (Hogerheijde & van der Tak 2000) using collisional cross sections from Dumouchel et al. (2010). The HNC abundance models follow the form in Qi et al. (2008), which assumes a molecular layer with a vertically constant abundance, bound by surface (\(\sigma_s\)) and midplane (\(\sigma_m\)) boundaries defined with respect to the vertically integrated hydrogen column in the unit of \(1.59 \times 10^{21}\) cm\(^{-2}\). We set the HNC boundaries using the values from Qi et al. (2008) for HCN (\(\log\sigma_s=-0.5\), \(\log\sigma_m=1.5\)). The
Graninger et al.

Figure 4. Comparison of HNC 3–2 data (upper left panel) and models with an outer ($R_o$) and hole ($R_h$) radius. The contours are $[3,4,5]\sigma$, based on the observed RMS noise. Images are generated using the EXT data only to display the small-scale emission structures.

The radial distribution of HNC was modeled using an abundance power-law with an outer radius ($R_o$), hole radius ($R_h$), and power-law index of -1.5. The parameter $R_o$ was varied between 35 and 60 AU and $R_h$ was varied between 20 and 50 AU. The minimum $\chi^2$ value was calculated from the modeled and observed HNC emission in the ($u$, $v$)-plane using the data from the compact and extended configuration.

Figure 4 shows examples of emission models and observations in the extended configuration. The best fit model based on the $\chi^2$ value is a 10 AU wide ring with a $R_o$ of 45 AU and $R_h$ of 35 AU (Figure 5), shown in the lower left panel of Figure 4. This model results in no significant residuals when subtracted from the data in the ($u$, $v$)-plane and then imaged. The observed HNC azimuthal asymmetry toward TW Hya is thus not significant. The presence of a HNC ring in the TW Hya disk is by contrast a statistically significant result. Figure 5 shows the models that fit the data within 2 and 3$\sigma$. Within 2$\sigma$, the hole size is 28–45 AU (3$\sigma$: 26–47 AU) and the outer edge is at 37–55 AU (3$\sigma$: <57 AU).

4. DISCUSSION

In this study, we detect and spatially resolve HNC emission toward TW Hya and HD 163296. From these observations we extracted disk averaged HNC/HCN emission ratios, and the inner and outer HNC emission radii in the TW Hya disk. Here, we will discuss the constraints that these observations provide on the HNC chemistry and the disk structure.

4.1. HNC/HCN Emission Ratios

The disk-averaged HNC/HCN emission ratio is 0.1 – 0.2 toward both TW Hya and HD 163296. Based on ISM observations (Schilke et al. 1992) and astrochemistry models (Graninger et al. 2014), the HNC/HCN ratio can be used to constrain the characteristic temperature of the emission region because of efficient destruction of HNC by hydrogen atoms at temperatures >25 K. However, to apply this result straightforwardly requires...
that HNC and HCN are co-spatial, that there is a simple relationship between emission and abundance ratios of HNC and HCN, and that temperature dominates the chemistry, none of which clearly hold in the case of these disks.

First, HNC and HCN are not co-spatial – some HCN is present toward the center of the disk, where HNC is absent. This may be partially mitigated by averaging the HNC and HCN emission only across disk areas where HCN and HNC emission are respectively detected. Second, the relationship between line emission and abundance ratios is complicated by high HCN line opacities and different HCN and HNC excitation properties (Sarrasin et al. 2010). We used the large velocity gradient radiative transfer code, RADEX (van der Tak et al. 2007), with the collisional rate constants from Dumouchel et al. (2011) to evaluate relationship between the HNC/HCN line emission and abundance ratios using temperatures and densities characteristic of HCN disk emission (Walsh et al. 2010) and the HCN column density from Qi et al. (2008). We find that the line and abundance ratios can differ by up to an order of magnitude, demonstrating that more detailed modeling is required to extract HNC/HCN abundance ratios from the observed line emission.

Even if abundance ratios are successfully extracted, it may be difficult to infer a characteristic emission region temperature. Other environmental characteristics than temperature may regulate the HNC/HCN abundance ratio in disks. The model by Graninger et al. (2014) neglects the influence of radiation, which is reasonable in cloud cores, but may not be a good approximation in disks (Meijerink et al. 2007). Loison et al. (2014) has also shown that the HNC/HCN ratio depends on the C/O ratio, which may be elevated in disks. In the absence of detailed disk chemistry models that incorporate the complete HNC/HCN chemistry, disk-averaged line ratios are thus not useful probes of disk properties.

4.2. The HNC inner and outer radii in the TW Hya disk

In the disk toward TW Hya, there is a lack of HNC emission interior to 28–45 AU, corresponding to temperatures above 21–13 K using the Qi et al. (2013b) temperature profile. HCN does not have an emission hole (Qi et al. 2008). The observed absence of HNC close to the star is consistent with temperature dependent HNC destruction (Graninger et al. 2014). It is difficult to explain the hole with the increased UV flux present at smaller radii, since HNC and HCN are thought to have similar photodissociation cross-sections (Roberge et al. 1991, van Dishoeck et al. 2006).

HNC also has an outer emission radius between 35–55 AU, with the best fit at 45 AU, substantially smaller than the observed radii of other gases, including CO, HCO$^+$ and HCN (Qi et al. 2008). This outer cutoff in HCN emission was unexpected. It coincides with the onset of C$_2$H emission, which is proposed to be a UV photochemistry product (Kastner et al. 2013), and is close to the radius where the dust density sharply decreases (Andrews et al. 2012). These spatial coincidences suggest three potential explanations for the outer cutoff of HNC emission: 1. UV photodissociation of HNC due to deeper penetration of stellar and interstellar UV photons in the low-column density outer disk, 2. an increased UV photodissociation of H$_2$ and carbon-bearing molecules accelerating the HNC$+H$ and HNC$+C$ destruction processes at a given temperature (Graninger et al. 2014, Loison et al. 2014), and 3. grain-gas thermal decoupling due to lower gas and grain densities, causing a temperature reversal in the outer disk and thus the increase in thermally controlled HNC destruction (i.e. a similar process that sets the inner edge).

The first two explanations both depend on a more efficient UV photochemistry exterior to 45 AU, consistent with the onset of C$_2$H emission at the same radius (Kastner et al. 2013), and signs of efficient UV photochemistry in the outer low-density regions of other disks (Oberg et al. 2011). Flared disks can intercept stellar UV radiation out to larger radii, and Lyo emission is efficiently down-scattered through disk towards the midplane (Fogel et al. 2011). Additionally, the interstellar radiation field will contribute a non-negligible amount of UV radiation to the outer disk (Bergin et al. 2003). Efficient UV penetration at large disk radii may increase the HNC photodissociation rate, reducing its abundance. This scenario could be tested by comparing HNC and HCN emission profiles. HNC and HCN are believed to have similar photodissociation cross-sections (Roberge et al. 1991, van Dishoeck et al. 2006), and if their outer emission radii are UV controlled, they should be the same. The best-fit outer HCN emission radius is ~100 AU (Qi et al. 2008), twice that of HNC, but this value is not well enough constrained to exclude direct photodissociation of HNC as an explanation for the HNC outer radius.

The second and third proposed explanations for the HNC outer cut-off should produce a smaller HNC than HCN emission radius, due to an increased HNC-to-HCN conversion rate in the outer disk. In the second scenario and increased UV penetration increases the atomic H and C abundance through H$_2$ and CO photodissociation.
Graninger et al. (Akimkin et al. 2013), speeding up HNC isomerization at a given temperature (Graninger et al. 2014). In the third scenario, a lower density in the outer disk reduces gas-grain thermal coupling sufficiently that direct gas heating by the interstellar radiation field becomes important. This heating could result in a radial temperature gradient reversal outside of 45 AU, again increasing the HNC isomerization rate.

Higher spatial resolution observations with ALMA should be able to distinguish between the three scenarios. The relative sizes of the HNC and HCN outer radii should differentiate between #1 on one hand and #2 and #3 on the other. If #1 is ruled out, the radial temperature profile of HCN beyond 45 AU should be able to differentiate between #2 and #3. Whichever scenario is confirmed, spatially resolved HNC/HCN emission modeled using non-LTE techniques clearly has a large potential to probe interesting disk properties.

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