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Temperature characteristics of epitaxially grown InAs quantum dot micro-disk lasers on silicon for on-chip light sources

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Temperature characteristics of optically pumped micro-disk lasers (MDLs) incorporating InAs quantum dot active regions are investigated for on-chip light sources. The InAs quantum dot MDLs were grown on V-groove patterned (001) silicon, fully compatible with the prevailing complementary metal oxide-semiconductor technology. By combining the high-quality whispering gallery modes and 3D confinement of injected carriers in quantum dot micro-disks, we achieved lasing operation from 10 K up to room temperature under continuous optical pumping. Temperature dependences of the threshold, lasing wavelength, slope efficiency, and mode linewidth are examined. An excellent characteristic temperature $T_o$ of 105 K has been extracted.

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Whispering gallery mode (WGM) micro-resonators incorporating quantum dot (QD) active regions have aroused great interest recently for on chip light sources, due to their low threshold, high temperature stability of operation, and small form factor.1–4 While suspended disks with air-clad waveguides provide the best optical confinement, the larger thermal resistance compared to ring resonators may be a drawback.3 Room temperature continuous-wave (CW) operation of InAs quantum dot micro-disk lasers (MDLs) has been demonstrated, but only limited research has been carried out on their temperature characteristics.2,4,5 Under CW excitation, Ide et al.3 reported MDL on a GaAs substrate with $T_o$ of 64 and 36 K in the temperature range of 130–230 K and above 230 K, respectively. Under pulsed excitation, Yang et al.5 demonstrated $T_o$ of 31 K above room temperature in MDL on a GaAs substrate. A systematic characterization of lasing behavior as a function of temperature for MDLs integrated on a silicon substrate is yet to be investigated.

Recently, we demonstrated InAs quantum dot micro-disk lasers grown on a V-groove patterned (001) Si substrate.5 In this work, we systematically characterized the lasing behavior of the MDLs, including threshold, lasing wavelength, slope efficiency, and mode linewidth, as a function of temperature. We have obtained an excellent characteristic temperature $T_o$ of 105 K for the epitaxially grown MDLs on Si, which outperform previously reported InAs quantum dot micro-disk lasers on GaAs substrates.5,5,7 These temperature insensitive microdisk lasers with compact size and low thresholds directly grown on (001) silicon can be attractive on-chip light sources to meet the low-power consumption and athermal performance demands of silicon photonic devices.

We first grew GaAs/Si templates consisting of 1 μm GaAs thin film on a V-groove patterned Si substrate by metal-organic vapor deposition (MOVCD), using the method detailed in Refs. 8 and 9. On the GaAs-on-Si template, microdisk laser structure was deposited in a molecular beam epitaxy (MBE) system. The MBE epi-layers consist of, from bottom to top, a 1 μm GaAs buffer, a 600 nm Al$_{0.7}$Ga$_{0.3}$As post region, and a 500 nm disk region where five-stacked InAs/GaAs dots-in-a-well (DWELL) structure is sandwiched by 50 nm Al$_{0.4}$Ga$_{0.6}$As cladding layers.10 Fig. 1(a) shows the photoluminescence (PL) emission spectra of the as-grown epitaxial structure, with pump power density varied from 20 to 4000 W/cm$^2$. Ground state emission (1.29 μm) is located within the O-band of fiber-optic communications (1.26–1.36 μm). The micro-disk lasers were defined using colloidal lithography, followed by BCl$_3$/Ar inductively coupled plasma (ICP) etching, as shown in Fig. 1(b). A dilute suspension of SiO$_2$ spheres in isopropyl alcohol (IPA) was dispersed onto the sample surface using a pipette and then rapidly dried on a hot plate at 110 °C. The commercial SiO$_2$ spheres are 4 μm in diameter with ±5% variation. Using the isolated sphere-masks, ICP etching was performed to reach a depth of ~1.1 μm into the epi-structure, stopping at the GaAs buffer layer. The disk membrane was suspended upon the virtual GaAs-on-Si substrate by partially etching the underlying Al$_{0.7}$Ga$_{0.3}$As sacrificial layer using dilute hydrofluoric (HF) acid. The lateral undercut was designed to be ~1.6 μm to guarantee efficient light confinement within the disk perimeter while maintaining good heat sinking through the underlying pedestals. Schematic illustration of the final device structure is presented in Fig. 1(c). Noteworthy is the unique defect trapping effect provided by growing GaAs on V-groove nano-patterned Si substrates, as indicated by the cross-sectional TEM image near the hetero-interface in Fig. 1(d). It allows us to...
achieve high quality III–V/Si epitaxy on exact (001) silicon substrates without incorporating intermediate Ge or other buffer layers.\textsuperscript{11} Better compatibility with the prevailing complementary metal oxide-semiconductor (CMOS) manufacturing processes and silicon-on-insulator (SOI) technologies differentiates this epitaxial growth platform with other quantum dot lasers grown on Ge-on-Si or offcut Si substrates.\textsuperscript{12–17}

The fabricated devices were optically pumped at normal incidence with a CW diode laser operating at 532 nm. The excitation beam was focused to approximately 4 \( \mu \text{m} \) in diameter. The laser emission was collected through the same objective and analyzed by a monochromator with an InGaAs charge-coupled device array. For temperature dependent measurements, the sample was mounted in a helium gas flow cryostat, with device temperature controlled from 10 K to room temperature. Fig. 2(a) compares the normalized lasing spectra at three times threshold pump power from 10 to 300 K. A scanning electron micrograph (SEM) image of an optically isolated, undercut MDL is shown in the inset of Fig. 2(a). The spectra show narrow lines corresponding to high-quality whispering-gallery modes. At 10 K, multimode lasing occurs. The mode at longer wavelengths wins in the mode competition and dominates the spectrum above 80 K. We observe red-shifting of the lasing peak with increasing temperature. Abrupt changes towards longer wavelength cavity modes were also observed due to the effect of the mode overlapping with the gain spectrum. Fig. 2(b) presents the integrated output-power intensity versus the pump power (I–I curve) of the lasing peak from 10 to 300 K. The distinct kinks signify the onset of laser operation over the entire temperature range.

Fig. 3(a) highlights the optical transition positions in the gain medium pumped by the lowest and highest excitation at 10 K. The emission below threshold (8.25 \( \mu \text{W} \)) was magnified 100\( \times \) and plotted together with that well above the threshold. A scanning electron micrograph (SEM) image of an optically isolated, undercut MDL is shown in the inset of Fig. 3(a). The spectra show narrow lines corresponding to high-quality whispering-gallery modes. At 10 K, multimode lasing occurs. The mode at longer wavelengths wins in the mode competition and dominates the spectrum above 80 K. We observe red-shifting of the lasing peak with increasing temperature. Abrupt changes towards longer wavelength cavity modes were also observed due to the effect of the mode overlapping with the gain spectrum. Fig. 2(b) presents the integrated output-power intensity versus the pump power (I–I curve) of the lasing peak from 10 to 300 K. The distinct kinks signify the onset of laser operation over the entire temperature range.

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threshold (495 μW). Four lasing peaks at 1162 nm, 1188 nm, 1191 nm, and 1220 nm were observed, evidenced by the pronounced kinks in the L–L curves in Fig. 3(b). The sub-peak (1188 nm) close to the main peak (1191 nm) arises from out-of-plane polarization. The other three well-separated modes correspond to WGM with different azimuthal orders, and the mode spacing is extracted to be around 29 nm. Using the equation of free spectral range (FSR) \( \Delta \lambda \approx \frac{k^2}{2\pi R} \), where R is the radius of the disk and \( \lambda \) is the emission wavelength, the effective group index is found to be 3.95. The relative slope efficiencies of the four modes are summarized in Fig. 3(c). The modes at longer wavelengths have larger slope efficiency compared to the ones at shorter wavelengths. This was accompanied by a moderate emission droop for the modes at shorter wavelengths (1188 nm and 1162 nm) with higher injection power, as shown by the L–L curves in Fig. 3(b). The higher slope efficiency of the longest wavelength performance suggests a higher capture efficiency in the QDs, as well as a slightly better confinement of the electron-hole pairs, leading to higher probability of radiative recombination.

At 300 K, overlay of the emission spectra at the two extreme pump powers (8.25 μW and 495 μW) in Fig. 4(a) shows that the mode on the low-energy side wins in the mode competition and dominates the spectrum. The winning of the low-energy modes in the mode competition is related to the larger QD size, corresponding to the lower energy emission, and the superior capture efficiency. Mode spacing at this temperature slightly increases to 35 nm compared to the value of 29 nm at 10 K, corresponding to a lower effective group index of 3.84 at 300 K.

A rate equation analysis was carried out to fit the experimental L–L curve to extract the spontaneous emission factor \( \beta \). We adopted a coupled rate equation model for the carrier density \( N \) and the photon density \( P \)

\[
\frac{dN}{dt} = \eta \frac{P_{in}}{\hbar \omega V} - V_g P B N^2 - C N^3, \tag{1}
\]

\[
\frac{dP}{dt} = V_g P + \beta B N^2 - \frac{P}{\tau_p}, \tag{2}
\]

where \( P_{in} \) is the pump power, \( \hbar \omega \) is the emitted photon energy, \( V_g \) is the group velocity (7.8 x 10^5 cm/s), B is the bimolecular recombination coefficient (2 x 10^{-10} cm^3/s), and C is the Auger coefficient (8 x 10^{-29} cm^3/s). The photon lifetime \( \tau_p \) (4.1 ps) is evaluated using the cold cavity quality factor (2950) obtained from sub-threshold measurements. Steady state solutions were found, and the gain function of \( g = \frac{g_0 N - N_0 N}{1 + \tau \tau_p} \), where \( g_0 \) is the differential gain (3 x 10^{-16} cm^2), \( \varepsilon \) is the gain compression factor (1 x 10^{-16} cm^2), and \( N_0 \) is the transparency carrier density (1.23 x 10^{18} cm^3). Fig. 4(b) presents the experimental data and rate equation analytical fits using different values of \( \beta \). The best fit to the experimental data reveals a high spontaneous emission coupling efficiency of 10%. The nonlinear kink in the log-log L–L curve also occurs at the same pump power values as the kink in the linear scale L–L version (inset in Fig. 4(b)).

Figs. 5(a) and 5(b) present the temperature dependence of the lasing wavelength and spectral linewidth at twice the threshold pump power, respectively. The spectral positions of the mode redshift slightly at a rate of \( \sim 0.04 \text{ nm/K} \) from 10 K to 80 K, and from 220 K to 280 K, which is ascribed to the temperature dependence of the cavity effective refractive index. The lasing line is at the same time governed by the temperature shrinkage of the active region bandgap, causing mode hopping towards longer wavelengths at 80 K.

FIG. 3. (a) Laser emission spectra measured below (yellow, 8.25 μW) and above (light blue, 495 μW) threshold in 10 K. Spectrum taken at 8.25 μW was amplified 100 times to be visible; (b) L–L curves of the lasing peak at 1162 nm, 1188 nm, 1191 nm, and 1220 nm in 10 K, and the dashed lines represent linear fits to the experimental data; and (c) threshold power as well as relative slope efficiency of the four modes at 1162 nm, 1188 nm, 1191 nm, and 1220 nm in 10 K.
and 280 K. The overall trend of the temperature dependent lasing line is governed by the confinement of the QDs, and the dependence of that confinement as a function of temperature, as well as the theoretical InAs band-gap shrinkage, as shown in Fig. 5(a). The theoretical InAs band-gap shrinkage is estimated via Varshni’s formula, \( \Delta E = \frac{A}{T^2} - \frac{B}{T} \), assuming \( A = 0.00042 \text{ eV/K}^2 \) and \( B = 199 \text{ K} \). In Fig. 5(b), the line-width slightly broadens to 0.4 nm. We assume this to be related to the thermal redshift during the accumulation time of the spectrum measurement, which is a more prominent factor at higher temperatures.\(^4\) There is a certain fluctuation of line-width with temperature, starting broad then reducing and then again increasing in irregular steps. This is probably due to varied amount of the spatial overlap of the QD gain spectral peak with the position of the cavity mode at different temperatures.

The threshold pump power as a function of temperature in Fig. 5(c) shows that the threshold increases by a factor of \( \sim 2 \) as the temperature increases from 10 K to 300 K. This can be fit with an exponential function using \( P_{th} \propto \exp(T/T_0) \),\(^19\) and the characteristic temperature \( T_0 \) was extracted to be around 105 K. A slight decrease in the threshold from 10 K to 80 K is due to the suppressed multimode lasing at 80 K. Otherwise, the experimental data follow the exponential fit, which increases at higher temperatures due to restrained carrier confinement in the QDs and enhanced non-radiative recombination. The relative slope efficiency decreases accordingly with increasing temperature, shown in Fig. 5(d). This is ascribed to the combined effects of free carrier absorption, intervalence band absorption, and suppressed constrains of defects.\(^7\)

Figs. 6(a) and 6(b) plot the histograms of the threshold over a number of devices at 10 K and 300 K, respectively. The average threshold is calculated to be 160 \( \mu \text{W} \) at 10 K and 250 \( \mu \text{W} \) at 300 K, as illustrated in Fig. 6(c). The overall device thresholds only increase by a factor of \( \sim 1.56 \) when the temperature ramps from liquid helium temperature (10 K) up to room temperature (300 K), suggesting good thermal stability.

In conclusion, we performed systematic characterization of lasing behavior of the monolithically integrated micro-disk lasers on a standard (001) silicon substrate. A high characteristic temperature \( T_0 \) of 105 K was extracted through the exponential fit of threshold power as a function of temperature from 10 K to 300 K. At 300 K, a high spontaneous emission coupling efficiency up to 10% is evidenced through rate equation analysis. The decent thermal stability and lasing characteristics suggest feasibility of manufacturing silicon chips with integrated micro-size lasers, promising to realize high-performance on chip optical links for telecommunication networks.
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