MEASUREMENTS OF T\textsubscript{0} TEMPERATURES OF SUPERSATURATED Si-As ALLOYS

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ABSTRACT

The congruent melting point, or \( T_0 \) curve, of crystalline Si-As alloys has been measured in the range of 16.6 to 18.1 at.\% arsenic by line source electron beam annealing. Alloys were created by ion implantation of As into 0.1 from Si-on-sapphire and crystallized by pulsed laser melting. \( T_0 \) temperatures decrease from 1673\,\textdegree \,K at 2.0 at.\% As to 1516\,\textdegree \,K at 18.1 at.\% As. The results of these measurements are significantly higher than the previous results of studies using pulsed laser melting techniques. Advantages of the e-beam technique over previous techniques are discussed. Chemical free energy functions of the solid and liquid phases were calculated from existing thermodynamic data. The calculated \( T_0 \) curve agrees with the measured values only in low concentration region (less than 8 at.\%).

INTRODUCTION

A quantitative experimental test of theories for nonequilibrium interface kinetics during alloy solidification [1] requires, as input, reliable thermodynamic functions for the solid and liquid phases. The Si-As system is a promising one for such a test because of the ease with which both of the solidification velocity and the nonequilibrium partition coefficient can be measured in Si alloys, and because of the high solubility of As in Si [2] and the large supersaturations that can be obtained in this system by rapid solidification [3]. Although As is one of the important dopants of Si, however, thermodynamic data in this system are insufficient and often incomplete [4-9]. One example is the \( T_0 \) curve, the temperature at which solid and liquid phases of equal composition have equal Gibbs free energies. Recently, pulsed laser melting experiments have been used to determine the \( T_0 \) temperatures of supersaturated Si-As alloys [8,9]. The \( T_0 \) temperatures reported in these experiments differ from each other by more than 200 degrees. Solid solubility data also contradict each other [4].

In order to obtain a useful thermodynamic constraint on the free energy functions, we re-measured the \( T_0 \) temperatures. Supersaturated Si-As alloy films on sapphire substrates were prepared by rapid melting and solidification with a pulsed excimer laser. Line source e-beam annealing (LEBA) was used for melting the alloy films [10-12]. The heating rate by e-beam irradiation is fast enough to suppress melting when partitioning is required, such as the solidus temperature is surpassed. Rather, melting initiates when the \( T_0 \) temperature is reached and continuous melting becomes possible. The LEBA experiment has several advantages over pulsed laser experiments. First, electrons are not reflected from the specimen surface, so accurate data on the reflectivity and on the reflectivity change upon melting are not required for a reliable calculation of the energy deposition. Second, the time scale of the e-beam heating experiment is on the order of 100\,\textmu\,s. The deposited energy can spread to greater depths in this amount of time. Therefore, the absorption depth is not important. Furthermore, the alloy films used in this work are thin enough (100\,\textmu\,m) that their effect on the e-beam energy deposition and thermal diffusion is negligible. For example, heat flow calculations indicate that the temperature difference across the alloy layer is less than 1 degree. Hence the temperature of the Si-As film can safely be assumed to equal that of sapphire substrate at the boundary. Consequently, the temperature history can be calculated by using only the physical data of sapphire, which are well known [10,11], and the e-beam parameters. These advantages make it possible to obtain more reliable \( T_0 \) temperatures in laser melting experiments [8,9].

EXPERIMENTAL RESULTS

Silicon films of 100 nm thickness on sapphire substrates were implanted with 60keV \( ^{75}\text{As}^+ \) ions at various doses from 1\,x\,10\textsuperscript{16} to 8\,x\,10\textsuperscript{17}cm\textsuperscript{-2}. Samples were covered by a mask during the
implantation to make alternating stripes of Si-As alloy and pure Si, as shown in Fig. 1. This sample geometry embedded a standard (pure S)  for each sample for e-beam power calibration. The implantation regions in these samples were all amorphous and had non-uniform As depth profiles. XeCl excimer laser pulses (40ns pulse duration, 308nm wavelength) were used to homogenize and crystallize the samples. The samples were melted from 5 to 15 times depending on the implant dose. The energy density of the laser pulse was selected to melt barely to the back of the film in order to avoid possible damage due to liquid overheating or long melt duration. Transient conductance and reflectance measurements were performed to record the melt history on a companion set of TCM-patterned, otherwise identical samples. After pulsed laser melting, the As concentration profiles were measured by Rutherford Backscattering Spectrometry (RBS) in grazing-exit geometry. The microstructure of one laser melted sample was observed in cross-sectional TEM to confirm that the resulting material is crystalline single phase, and to examine the density of crystalline defects. Low dose samples (1x10^16 and 2x10^16/cm^2) were crystallized and homogenized by melting and regrowth during the laser irradiation. It was observed that after the laser irradiation, the electric conductance increases and the color of the sample surface changes due to crystallization. However, high dose samples (7x10^16/cm^2 and 8x10^16/cm^2) regrew as amorphous after full melting. These samples were homogenized by a number of melting and regrowth cycles, and then explosively crystallized at low energy densities. Cross-sectional TEM of one explosively crystallized sample (dose: 6x10^16/cm^2) revealed single phase material with fine grains of about 30nm diameter, as seen in previous studies on explosive crystallization in pure Si. Samples with the intermediate dose of 4x10^16/cm^2 seemed to solidly partly crystalline and partly amorphous as judged by optical inspection. We subsequently used the same explosive crystallization technique as in high dose samples to fully crystallize these samples. The compositions of As in the homogenized samples are shown in Table 1. The final composition-depth profiles never displayed nonuniformities greater than ±2 at.%. A line source e-beam (20nm long x 1mm wide) was used to heat these samples rapidly to a known peak temperature. The e-beam source and annealing system are described in detail elsewhere [10-12]. Since the thermal properties of sapphire and the e-beam parameters are accurately known or measured, the temperature profile along the e-beam sweep can be calculated to within 10 degrees. The e-beam was swept along the implantation stripe as shown in Fig. 1. Hence, the e-beam could create Si-As alloy stripes and pure Si stripes simultaneously. The microstructure of the e-beam melted region is distinctly different from that of the unmelted regions. The e-beam power used in the temperature profile calculation was calibrated using the melt width in the pure Si stripes. T_e temperature was then determined by comparing the melt width of the alloy stripe with the calculated temperature profile. Measured T_e temperatures of supersaturated Si-As alloys are shown in column 3 of Table 1. Low dose samples (1x10^16 and 2x10^16/cm^2) were crystallized directly by regrowth from the melt. These samples are thus expected to have large grains, even though they contain some defects. The measured T_e values are thus highly reliable in the low As concentration range. However, since the high dose samples (2x10^16/cm^2 and 3x10^16/cm^2) were crystallized by explosive crystallization resulting in a fine polycrystalline microstructure, the T_e values from these samples remain preliminary until the effect of the microstructure on the T_e value is investigated. We have found that, depending on laser conditions, the 2x10^16/cm^2 samples can be either solidified in the crystalline state or solidified in the amorphous state and subsequently explosively crystallized. The T_e temperatures of these alloys will be compared to each other in order to determine whether there is any effect on the apparent T_e values due to the fine-grained microstructure that results from explosive crystallization.

**DISCUSSION**

The T_e temperatures measured in this work are summarized in Fig. 2 as a function of As concentration. Results from the pulsed laser melting experiments are also shown in Fig. 2. Baeri et al. [8] found much lower T_e temperatures than the work. They measured the threshold energy densities for melting by a pulsed ruby laser and then calculated the specimen surface temperature. This calculation required optical and thermal parameters such as reflectivity, absorption coefficient, heat capacity, and thermal diffusivity. For laser melting in the nanosecond time regime, the details of laser energy deposition and temperature gradients in the surface alloy layer are very important for calculations of the temperature profile. However, in the absence of thermophysical data, assumptions had to be made about the dependencies of these parameters on temperature and solute concentration. In the other pulsed laser experiment, Peery et al. estimated the T_e temperature from the retarding of solidification near the As peak in As-implanted Si on sapphire samples [9]; they also found lower T_e temperatures than ours. Although the laser energy deposition did not affect their results, assumptions had to be made about the composition dependence of thermal properties of the alloy in their calculation. The uncertainties in the thermal and optical properties of Si-As alloys should be considered the main reason for the inaccurate estimation of T_e values in both previous studies.

This measurement of the T_e curve can be used as a constraint on thermodynamic modeling of the solid and liquid phases in these alloys. We derived the free energy functions of each phase from existing thermodynamic data. The validity of the thermodynamic model will be tested by comparing calculated T_e curves with the measured values. We assumed a quasi-equilibrium model (the interaction parameter was allowed to depend linearly on temperature, but not on composition) for both the solid and liquid phases. With the diamond cubic phases of pure Si and pure As chosen as the reference states, the molar Gibbs free energy for each phase can be written

\[ G^f = RT \left[ (1-x) \ln (1-x) + x \ln (x) \right] + x \Omega_1 \]

\[ G^l = (1-x) \Delta \Delta G_{Si(d=0)} + x H(\theta) \]

\[ + RT \left[ (1-x) \ln (1-x) + x \ln (x) \right] + x \Omega_1 \]

\[ \Omega_2 \]
Here, $x_j$ is the mole fraction of As in phase $i$, $R$ the gas constant, and $T$ the temperature in K. The "lattice stability" of pure Si, $\Delta G^o_{\text{Si(Si)}}$, is known [14]. The unknown "lattice stability" of pure diamond cubic As, $\Delta G^o_{\text{As(As)}}$, was estimated by assuming that the melting temperature of pure diamond cubic As would be in the range from 0K to -5000K, and that the energy of fusion is the same as for pure Si. The thermodynamic calculations reported here were done using both 0K and -5000K for the melting temperature of diamond cubic As to provide reasonable estimates of upper and lower limits on the calculated $T_0$ values. The interaction parameters of each phase, $\Omega_i$ and $\Omega_j$, were obtained by fitting to the activity coefficients of As in Si:As solid solutions [15,16] and to the experimental liquidus data [4,17], respectively. The thermodynamic parameters obtained as linear functions of temperature are listed in Table 2.

The solid and broken lines in Fig.2 are the calculated phase boundaries and $T_0$ curves when 0K and -5000K, respectively, are used as the melting temperature of diamond cubic As. The calculated phase boundaries and $T_0$ curves in the low concentration region are not sensitive to the uncertainty in the As lattice stability. However, the $T_0$ curves become more sensitive to such uncertainty at higher As concentrations. The calculated $T_0$ curve agrees with our measurements for As concentrations less than 8 at.%. Beyond this composition, the calculated $T_0$ temperature decreases faster than the measured $T_0$ temperatures. Although the measured values in the high concentration region are preliminary, it is anticipated that any "microstructure correction" would only raise these data points to higher values, so it is unlikely that agreement with the model will result. One possible reason for the deviation in the high As concentration region is that the available thermodynamic data are insufficient. For example, the activity coefficient of As in Si:As solid solutions was measured only at low concentrations (maximum 3 at. % As) [15,16].

### Table 2. Calculated thermodynamic parameters of Si-As alloy in Jmol.

This data is based on extrapolation from the high concentration range. The model used in this work is too simple to consider the thermodynamic behavior in the high concentration regime. However, a model with more free parameters also requires more thermodynamic data to determine them.

### CONCLUSIONS

The combination of rapid quenching by pulsed laser melting and sufficiently rapid heating by electron-beam irradiation provide a useful technique to measure the $T_0$ temperature of supercooled alloys. The use of the e-beam melting technique with thin film samples has the important advantage that the thermophysical properties of the alloy layer and the e-beam parameters are not important for the temperature calculation. We observe higher $T_0$ temperatures for the Si-As system than do previous experiments [8,9].

The $T_0$ curve was calculated using available thermodynamic data for dilute alloys. This calculation agrees with the measured data in the low As concentration region. However, the calculations predict lower $T_0$ temperatures that are measured for high As concentrations. Further studies are in progress to obtain a thermodynamic model valid over the entire composition range.

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### REFERENCES

NEW INSIGHT INTO THE BANDED MICROSTRUCTURE OF RAPIDLY SOLIDIFIED Ag-Cu ALLOYS: EXPERIMENT AND THEORY

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ABSTRACT
A transmission electron microscopy study has revealed that the banded microstructure produced in Cu-61.7 at. pct. Ag after surface laser metal and resolidification is not due to a cellular breakdown of the interface as previously thought, but consists of domains or regions of the extended metastable solid solution, y, and a coupled growth structure containing thin plates of y and the copper-rich phase \( \beta' \). The spacing of this coupled growth structure is approximately 100 \( \AA \), which is less than half the minimum spacing yet observed for coupled growth from the melt in Ag-Cu.

In order to explain the banded microstructure a theoretical model has been developed based on finite elements solution of the diffusion equation in the melt, the interface response functions for continuous growth, and the banded growth developed by Aziz and Kaplan, and a recently developed theory for coupled growth that includes far-from-equilibrium regimes by Yankov et al.

INTRODUCTION
Although the banded structure that appears in Ag-Cu alloys after laser beam surface melting and resolidification has been known for some time, its microscopic manifestation or growth theory has been developed to explain this phenomenon. As revealed by electron microscopy, the banded microstructure consists of equilaxed spaced grooves, marking sequential positions of the solid-liquid interface during solidification, with a typical spacing of about 1 \( \mu \text{m} \). Since it was first observed about 20 years ago in pulsed laser melting of Ag-Cu alloys [1], the banded microstructure has presented a challenge to researchers as it did not correspond to any of the solidification patterns that have been observed and modeled until now.

Two major experimental studies did not satisfactorily explain the origin of the bands. Beck, Coppley and Bas (2.3) melted a thin layer on the surface of Ag-Cu specimens of different compositions with a scanning CW CO\(_2\) laser and subsequently performed scanning electron microscopy (SEM) and X-ray diffraction studies of the banded region. They suggested that the bands may originate by a diffusional instability described by Baker and Cahn (4). They also established using energy dispersive spectroscopy a mean increase of about 1.1 at. pct. Ag in the bands. Boettinger et al. (5) performed similar experiments substituting an electron beam for the laser. While this work, at higher scanning speeds (100-200 cm s\(^{-1}\)), transition between the banded structure and disordered solidification was observed. On the other hand the maximum speed at which coupled growth was found to exist was on the order of 2.5 cm s\(^{-1}\) with a spacing of about 20 nm. Based on transmission electron microscopy observations of a cellular-like structure at one non-eutectic composition, they suggested that the bands may be due to a periodical cellular breakdown of the solidification front near the limit of absolute stability (8). However microscopic evidence was not presented for the other compositions they investigated in order to be able to demonstrate this with certainty. A modified analysis of the original stability theory proposed by Coriell and Sakeha (7) predicts that the partition coefficient, \( \chi \), varies with velocity and the derivative \( \chi \text{d} \chi \text{d}v \text{d} \) is large, oscillatory instabilities may form. However, these would produce a periodic solution variance in the growth direction, which were not observed by Boettinger et al.

In this study we present transmission electron microscopy (TEM) evidence for the microstructure of the banded structure in an alloy of near eutectic composition. The results show that the bands consist of alternating regions of the extended solid solution and a coupled growth structure. The highlights of a model developed to explain the periodic appearance of these metastable structures is also presented.

EXPERIMENTAL PROCEDURES AND RESULTS
The samples were prepared from 99.99 at. pct. pure Cu and 99.999 at. pct. pure Ag to produce an alloy...